

CKP Hierarchy, Bosonic Tau Function and Bosonization Formulae^{*}

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Abstract. We develop the theory of CKP hierarchy introduced in the papers of Kyoto school [Date E., Jimbo M., Kashiwara M., Miwa T., *J. Phys. Soc. Japan* **50** (1981), 3806–3812] (see also [Kac V.G., van de Leur J.W., *Adv. Ser. Math. Phys.*, Vol. 7, World Sci. Publ., Teaneck, NJ, 1989, 369–406]). We present appropriate bosonization formulae. We show that in the context of the CKP theory certain orthogonal polynomials appear. These polynomials are polynomial both in even and odd (in Grassmannian sense) variables.

Key words: integrable system; Pfaffian; Hafnian; symmetric functions; Schur type functions

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1 Introduction

In this paper we develop the ideas of Date, Jimbo, Kashiwara and Miwa [2, 3]. In [3] it was pointed out that the tau function of the CKP hierarchy may be presented as the vacuum expectation value of bosonic fields ϕ_n which act in a bosonic Fock space, denoted by F in the present paper. As it was shown in [3] the higher CKP flows are induced by the action of bosonic current algebra operators J_n , $n > 0$. We shall show that in contrast to the familiar fermionic approach the action of the currents J_{-n} on the vacuum state does not generate the whole Fock space F where the original bosonic operators ϕ_n act. To generate the whole Fock space we need to add an additional fermionic field which is a sort of super-counterpart to the bosonic current. This problem is studied in Section 2. In Subsection 2.1 we introduce Fermi operators whose action on the vacuum vector complete the action of the current algebra J_n to obtain the whole Fock space F . A bosonization formula which expresses the original bosonic field $\phi(z)$ in terms of the current algebra J_n and of the fermion field $\theta(z)$ is suggested (2.11). (As a byproduct of this relation we obtain an equality (2.29) which relates Pfaffian and Hafnian expressions which earlier appeared in a quite different context [7].) Here we show that this fermionic field is a super-partner of currents and it naturally creates a dependence of CKP tau function on auxiliary odd Grassmannian parameters. Though we present a bilinear equations written in terms of super vertex operator we do not construct Lax equations with respect to odd parameters. In Section 4 we introduce new orthogonal polynomials in many variables, $C_\lambda(\mathbf{t})$, which appear as a result of

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the ‘bosonization’ of the basis Fock vectors of F , see (4.4). These polynomials depend both on CKP higher times t_n (n odd) and the above-mentioned Grassmannian odd parameters. In certain sense these polynomials play a role similar to the role of the Schur functions in the theory of KP [16] and TL [17], and the role of the projective Schur functions in the theory of BKP [12, 19], namely, CKP tau function may be presented as a series in these polynomials over partitions (see Subsection 4.6). However in contrast to the KP and BKP cases these polynomials are not CKP tau functions themselves. At the end some combinatorial properties of C_λ are discussed.

2 CKP bosonic tau function

In this section we follow a suggestion in [3] and describe a CKP hierarchy of PDEs starting from a collection of free bosons. This imitates their approach in the BKP case, where one starts with neutral fermions. However, this hierarchy, although related to Lie algebra c_∞ , differs from the usual CKP hierarchy, for which one takes a reduction of the KP hierarchy by assuming that the Lax operator satisfies $L^* = -L$; such Lax operators come from certain KP tau functions, which are fixed by some involution and where one puts the even times to zero, see e.g., [3] or [1] for more details. The hierarchy described in this paper is different and is not related to the usual CKP which comes from a reduction of KP. In the latter case one has a realization of c_∞ , for which the level is positive. Our construction realizes c_∞ with a *negative* level.

2.1 Bose–Fermi correspondence in the CKP case

We follow a suggestion of Date, Jimbo, Kashiwara and Miwa in their paper [3] and introduce free bosons, but for convenience of notation we shift the index by $\frac{1}{2}$. So ϕ_i with $i \in \frac{1}{2} + \mathbb{Z}$ satisfy commutation relations:

$$\phi_i \phi_j - \phi_j \phi_i = (-)^{j-\frac{1}{2}} \delta_{i,-j}. \quad (2.1)$$

The Fock space F , respectively F^* , is defined by

$$\phi_j |0\rangle = 0 \quad \text{if } j < 0, \quad \text{resp.} \quad \langle 0 | \phi_j = 0 \quad \text{if } j > 0, \quad (2.2)$$

so that F has as basis the vectors

$$(\phi_{j_1})^{m_1} (\phi_{j_2})^{m_2} \dots (\phi_{j_{n-1}})^{m_{n-1}} (\phi_{j_n})^{m_n} |0\rangle \quad (2.3)$$

with $j_1 > j_2 > \dots > j_{n-1} > j_n > 0$ and m_i positive integers. Defining

$$\deg |0\rangle = 0, \quad \deg \phi_j = j,$$

we have a direct sum decomposition of F :

$$F = \bigoplus_{k \in \frac{1}{2}\mathbb{Z}} F_k \quad \text{with} \quad F_k = \{f \in F \mid \deg f = k\}.$$

It is straightforward to check that the dimension of F_k is given by the partition of k into positive elements of $\frac{1}{2} + \mathbb{Z}$. Define the formal character as

$$\dim_q F = \sum_{k \in \frac{1}{2}\mathbb{Z}} \dim F_k q^k.$$

Then

$$\dim_q F = \prod_{0 < k \in \frac{1}{2} + \mathbb{Z}} \frac{1}{1 - q^k}. \quad (2.4)$$

Writing

$$\phi(z) = \sum_{j \in \frac{1}{2} + \mathbb{Z}} \phi_j z^{j - \frac{1}{2}},$$

we denote

$$H(z) := \sum_{n \in 1 + 2\mathbb{Z}} J_n z^{-n-1} := -\frac{1}{2} : \phi(-z) \phi(z) :, \quad (2.5)$$

where the normal ordering is defined by

$$:\phi_i \phi_j: = \begin{cases} \phi_i \phi_j & \text{if } i \geq j, \\ \phi_j \phi_i & \text{if } j > i. \end{cases} \quad (2.6)$$

In other words $J_n = 0$ for n even and

$$J_n = \frac{1}{2} \sum_{j \in \frac{1}{2} + \mathbb{Z}} (-)^{j + \frac{1}{2}} \phi_j \phi_{-j-n} \quad \text{for } n \text{ odd};$$

one has the following familiar commutation relations

$$[J_n, J_m] = -\frac{n}{2} \delta_{m, -n}. \quad (2.7)$$

The elements $:\phi_i \phi_j:$ form a representation of the Lie algebra c_∞ , see e.g. [9]. However, we want to stress that its level (the value of its central element) is negative. Note also that in the commutation relations (2.7) we have the factor $-\frac{n}{2}$ instead of the usual $\frac{n}{2}$.

It is clear that

$$J_n |0\rangle = \langle 0 | J_{-n} = 0 \quad \text{for } n > 0.$$

By a similar argument as before, again since these are bosons, we can apply an element J_{-n} infinitely many times to $|0\rangle$. Since the degree of J_n is $-n$ we obtain that the action of this Heisenberg algebra on the vacuum vector produces in the $\dim_q F$ the partition function of partitions in only odd numbers:

$$\prod_{0 < k \in 1 + 2\mathbb{Z}} \frac{1}{1 - q^k}.$$

Now we calculate, using (2.4),

$$\left(\prod_{0 < k \in 1 + 2\mathbb{Z}} \frac{1}{1 - q^k} \right)^{-1} \dim_q F = \prod_{0 < k \in \frac{1}{2} + \mathbb{Z}} \frac{1 - q^{2k}}{1 - q^k} = \prod_{0 < k \in \frac{1}{2} + \mathbb{Z}} (1 + q^k). \quad (2.8)$$

This part should be explained by something else and we expect it to be fermions, at least anticommuting variables. The factor $1 + q^k$ is related to a fermion of degree k . This is how we get these elements and calculate their commutation relations.

We first calculate

$$[J_n, \phi(z)] = \frac{1}{2} \sum_{j, k \in \frac{1}{2} + \mathbb{Z}} (-)^{j + \frac{1}{2}} [\phi_j \phi_{-j-n}, \phi_k] z^{k - \frac{1}{2}}$$

$$\begin{aligned}
&= \frac{1}{2} \sum_{j,k \in \frac{1}{2} + \mathbb{Z}} (-)^{j+\frac{1}{2}} ([\phi_j, \phi_k] \phi_{-j-n} + \phi_j [\phi_{-j-n}, \phi_k]) z^{k-\frac{1}{2}} \\
&= \frac{1}{2} \sum_{j,k \in \frac{1}{2} + \mathbb{Z}} (-)^{j+\frac{1}{2}} (-)^{k-\frac{1}{2}} (\delta_{j,-k} \phi_{-j-n} + \delta_{j+n,k} \phi_j) z^{k-\frac{1}{2}} \\
&= \frac{1}{2} \sum_{k \in \frac{1}{2} + \mathbb{Z}} 2\phi_{k-n} z^{k-\frac{1}{2}} = z^n \phi(z).
\end{aligned} \tag{2.9}$$

Now, using (2.7) we see that

$$[J_n, e^{\frac{2}{m} J_m z^{-m}}] = \delta_{n,-m} z^n e^{\frac{2}{m} J_m z^{-m}}. \tag{2.10}$$

Hence setting

$$\theta(z) := V_-(z)^{-1} \phi(z) V_+(z)^{-1}, \tag{2.11}$$

where

$$V_{\pm}(z) = \exp \sum_{\pm k > 0, \text{ odd}} \frac{2}{k} J_k z^{-k}, \tag{2.12}$$

we have, from (2.9) and (2.10),

$$[J_n, \theta(z)] = 0. \tag{2.13}$$

2.2 Commutation relations of the $\theta(z)$'s

We now want to calculate the commutation relations of these $\theta(z)$'s given in (2.11). For this we first rewrite the commutation relations (2.1) as follows:

$$\phi(z)\phi(y) - \phi(y)\phi(z) = \delta(z - (-y)),$$

where

$$\delta(z - y) = z^{-1} \sum_{k \in \mathbb{Z}} \left(\frac{z}{y} \right)^k.$$

Note also that

$$\phi(-z)\phi(y) = : \phi(-z)\phi(y) : - \frac{1}{z} \frac{1}{1 - \frac{y}{z}}.$$

We first show the following identities

$$\begin{aligned}
V_+(-z)^{-1} V_-(y)^{-1} &= \frac{1 - \frac{y}{z}}{1 + \frac{y}{z}} V_-(y)^{-1} V_+(-z)^{-1}, \\
\phi(-z) V_-(y)^{-1} &= \frac{1 + \frac{y}{z}}{1 - \frac{y}{z}} V_-(y)^{-1} \phi(-z), \\
V_+(-z)^{-1} \phi(y) &= \frac{1 + \frac{y}{z}}{1 - \frac{y}{z}} \phi(y) V_+(-z)^{-1}.
\end{aligned} \tag{2.14}$$

Introduce

$$V(t) = \exp \sum_{k > 0, \text{ odd}} t_k J_k.$$

Then

$$V_+(z) = V\left(\frac{2}{1}z^{-1}, \frac{2}{3}z^{-3}, \frac{2}{5}z^{-5}, \dots\right). \quad (2.15)$$

The first equation of (2.14) is obtained in the following way. First using (2.7) one has

$$\begin{aligned} V(t)V_-(y) &= \exp\left[\sum_{k>0, \text{ odd}} t_k J_k, -\sum_{\ell>0, \text{ odd}} \frac{2}{\ell} y^\ell J_{-\ell}\right] V_-(y)V(t) \\ &= \exp\left(\sum_{k>0, \text{ odd}} t_k y^k\right) V_-(y)V(t). \end{aligned}$$

Combining this with (2.15) one obtains

$$\begin{aligned} V_+(-z)^{-1}V_-(y)^{-1} &= \exp\left(-\sum_{k>0, \text{ odd}} \frac{2}{k} \left(\frac{y}{z}\right)^k\right) V_-(y)^{-1}V_+(-z)^{-1} \\ &= \frac{1 - \frac{y}{z}}{1 + \frac{y}{z}} V_-(y)^{-1}V_+(-z)^{-1}. \end{aligned}$$

Using $\phi(-z)J_n = (J_n - (-z)^n)\phi(-z)$, see (2.9), we obtain the second relation in (2.14) as follows:

$$\begin{aligned} \phi(-z)V_-(y)^{-1} &= \phi(-z) \exp\left(\sum_{k>0, \text{ odd}} \frac{2}{k} J_{-k} y^k\right) = \exp\left(\sum_{k>0, \text{ odd}} \frac{2}{k} (J_{-k} + z^{-k}) y^k\right) \phi(-z) \\ &= V_-(y)^{-1} \exp\left(\sum_{k>0, \text{ odd}} \frac{2}{k} \left(\frac{y}{z}\right)^k\right) \phi(-z) = \frac{1 + \frac{y}{z}}{1 - \frac{y}{z}} V_-(y)^{-1} \phi(-z). \end{aligned}$$

The third formula is proved in a similar way.

We will also use the following identities which can be found in V. Kac's book [8]:

$$\begin{aligned} (z-y)\partial_y \delta(z-y) &= \delta(z-y), \\ (z-y)^{k+1} \partial_y^k \delta(z-y) &= 0, \\ \delta(z-y)a(z) &= \delta(z-y)a(y), \\ \partial_y \delta(z-y)a(z) &= \partial_y \delta(z-y) (a(y) + (z-y)\partial_y a(y)). \end{aligned}$$

We now calculate

$$\begin{aligned} \theta(-z)\theta(y) &= V_+(-z)^{-1}\phi(-z)V_+(-z)^{-1}V_-(y)^{-1}\phi(y)V_+(y)^{-1} \\ &= \frac{1 - \frac{y}{z}}{1 + \frac{y}{z}} V_+(-z)^{-1}\phi(-z)V_-(y)^{-1}V_+(-z)^{-1}\phi(y)V_+(y)^{-1} \\ &= V_+(-z)^{-1}V_-(y)^{-1}\phi(-z)V_+(-z)^{-1}\phi(y)V_+(y)^{-1} \\ &= \frac{1 + \frac{y}{z}}{1 - \frac{y}{z}} V_+(-z)^{-1}V_-(y)^{-1}\phi(-z)\phi(y)V_+(-z)^{-1}V_+(y)^{-1} \\ &= \frac{1 + \frac{y}{z}}{1 - \frac{y}{z}} V_+(-z)^{-1}V_-(y)^{-1} \left(: \phi(-z)\phi(y) : - \frac{1}{z} \frac{1}{1 - \frac{y}{z}} \right) V_+(-z)^{-1}V_+(y)^{-1}. \end{aligned}$$

Now replacing z and y by $-y$ and $-z$ respectively, gives

$$\theta(y)\theta(-z) = \frac{1 + \frac{z}{y}}{1 - \frac{z}{y}} V_-(-z)^{-1} V_-(y)^{-1} \left(: \phi(-z) \phi(y) : + \frac{1}{y} \frac{1}{1 - \frac{z}{y}} \right) V_+(-z)^{-1} V_+(y)^{-1}$$

and thus

$$\begin{aligned} \theta(-z)\theta(y) + \theta(y)\theta(-z) &= 2z\delta(z-y)V_-(-z)^{-1}V_-(y)^{-1} : \phi(-z)\phi(y) : V_+(-z)^{-1}V_+(y)^{-1} \\ &\quad + \left(\frac{1}{y} \frac{1 + \frac{z}{y}}{(1 - \frac{z}{y})^2} - \frac{1}{z} \frac{1 + \frac{y}{z}}{(1 - \frac{y}{z})^2} \right) V_-(-z)^{-1}V_-(y)^{-1}V_+(-z)^{-1}V_+(y)^{-1} \\ &= 4yH(y)\delta(z-y) - \partial_y\delta(z-y)(y+z)V_-(-z)^{-1}V_-(y)^{-1}V_+(-z)^{-1}V_+(y)^{-1} \\ &= 4yH(y)\delta(z-y) - 2y\partial_y\delta(z-y) - (z-y)\partial_y\delta(z-y) \\ &\quad \times (1 + 2y\partial_y(V_-(-y)^{-1})V_-(y)^{-1} + 2y\partial_y(V_+(-y)^{-1})V_+(y)^{-1}) \\ &= 4yH(y)\delta(z-y) - 2y\partial_y\delta(z-y) - \delta(z-y)(1 + 4yH(y)) \\ &= -2y\partial_y\delta(z-y) - \delta(z-y) = -D_y\delta(z-y), \end{aligned} \tag{2.16}$$

where $H(y)$ is as in (2.5) (not the one in (2.28)), and $D_y = y\partial_y + \partial_y y$ is the Euler operator.

Now write

$$\theta(z) = 2 \sum_{i \in \frac{1}{2} + \mathbb{Z}} J_i z^{-i - \frac{1}{2}}. \tag{2.17}$$

Note that there is no conflict with the J 's defined in (2.5), since here the J 's have indices in $\frac{1}{2} + \mathbb{Z}$. It is clear that the above commutation relation (2.16) in modes gives

$$J_j J_k + J_k J_j = (-)^{j - \frac{1}{2}} \frac{j}{2} \delta_{j, -k}, \quad j, k \in \frac{1}{2} + \mathbb{Z}$$

(compare with (2.7)). From (2.13) we also have

$$[J_n, J_m] = 0, \quad n \in 1 + 2\mathbb{Z}, \quad m \in \frac{1}{2} + \mathbb{Z}.$$

Thus we can combine the (anti)commutation relations of all J 's as follows:

$$[J_i, J_j]_s = \frac{j}{2} (-1)^{[j - \frac{1}{2}]} \delta_{i, -j}, \tag{2.18}$$

where the notation $[\ , \]_s$ serves for the supercommutator while $[i]$ denotes the integer part of a real number i . As we see, $\deg J_i = -i$ and that

$$J_k|0\rangle = \langle 0|J_{-k} = 0, \quad J_{-k}|0\rangle \neq 0 \neq \langle 0|J_k \quad \text{for } k > 0.$$

2.3 Even and odd times

Since the elements J_{-k} with $k \in \frac{1}{2} + \mathbb{Z}$ anticommute among themselves, they can only appear once in

$$J_{-k_n} J_{-k_{n-1}} \cdots J_{-k_3} J_{-k_2} J_{-k_1} |0\rangle.$$

Such a J_{-k} explains the factor $1 + q^k$ in the q -dimension formula (2.8). One can identify the J_n 's, for $n < 0$ with even and odd times, i.e., with commuting variables t_j , $0 < j \in 1 + 2\mathbb{Z}$, and Grassmann variables $t_{\frac{j}{2}}$, $0 < j \in 1 + 2\mathbb{Z}$, and identify Fock space F with the space

$$\mathbb{C}[t_{2j-1}, t_{\frac{2j-1}{2}}; j = 1, 2, \dots]$$

(or some completion of it, since we take exponentials), where one has

$$t_i t_j - (-)^{4ij} t_j t_i = 0,$$

in particular, $t_j^2 = 0$ for $j \in \frac{1}{2} + \mathbb{Z}$. We will write $\mathbf{t} = (t_1, t_{\frac{1}{2}}; t_3, t_{\frac{3}{2}}; t_5, t_{\frac{5}{2}}; \dots)$ and use $t = (t_1, t_3, t_5, \dots)$ and $t_{\text{odd}} = (t_{\frac{1}{2}}, t_{\frac{3}{2}}, t_{\frac{5}{2}}, \dots)$.

Let σ be this isomorphism, sending F to $\mathbb{C}[t_{2j-1}, t_{\frac{2j-1}{2}}; j = 1, 2, \dots]$. Then

$$\sigma J_{-j} \sigma^{-1} = (-)^{[\frac{1}{2}-j]} \frac{j}{2} t_j \quad \text{and} \quad \sigma J_j \sigma^{-1} = \frac{\partial}{\partial t_j}, \quad j > 0, \quad (2.19)$$

give the field exactly in commuting and anticommuting variables t_k .

Now using the free boson-(boson+fermion) correspondence, i.e., using the vertex operator expressions for the fields

$$\begin{aligned} \sigma \phi(z) \sigma^{-1} &= \exp \left(\sum_{\substack{0 < k \in \mathbb{Z} \\ k: \text{odd}}} t_k z^k \right) \exp \left(\sum_{\substack{0 < k \in \mathbb{Z} \\ k: \text{odd}}} \frac{2}{k} \frac{\partial}{\partial t_k} z^{-k} \right) \\ &\quad \times \sum_{0 < j \in \mathbb{Z}} \left((2j-1) t_{\frac{2j-1}{2}} (-z)^{j-1} + 2 \frac{\partial}{\partial t_{\frac{2j-1}{2}}} z^{-j} \right) \end{aligned} \quad (2.20)$$

(where we used (2.12), (2.11), (2.17) and (2.19)), in the following subsections we shall express the bilinear identity as a hierarchy of differential equations. A similar expression for (2.20) was also found in [9].

2.4 The CKP bilinear equation

Following [3] we define the operator

$$S = \sum_{k \in \frac{1}{2} + \mathbb{Z}} (-)^{k+\frac{1}{2}} \phi_k \otimes \phi_{-k} \equiv \oint \phi(z) \otimes \phi(-z) \frac{dz}{2\pi i},$$

that S commutes with the action of $:\phi_i \phi_j:$ on the tensor product $F \otimes F$ of the Fock space F and

$$S(|0\rangle \otimes |0\rangle) = \sum_{k \in \frac{1}{2} + \mathbb{Z}} (-)^{k+\frac{1}{2}} \phi_k |0\rangle \otimes \phi_{-k} |0\rangle = 0.$$

The CKP Hirota equation is [3]:

$$\oint \phi(z) g |0\rangle \otimes \phi(-z) g |0\rangle dz = 0, \quad (2.21)$$

where g is for instance given by (see [3]):

$$g = \exp \left(- \sum_{m,n} c_{nm} : \phi_{n+\frac{1}{2}} \phi_{m+\frac{1}{2}} : \right). \quad (2.22)$$

We rewrite (2.21) as

$$\text{Res}_z \phi(z) g |0\rangle \otimes \phi(-z) g |0\rangle = 0. \quad (2.23)$$

We could now use the isomorphism σ to define this hierarchy in terms of the times \mathbf{t} . However we will not do that yet, but concentrate first in the next subsection on the form of $\sigma(g|0)$.

2.5 The CKP tau function

Now let ODP_{ev} be the set of all partitions in an even number of odd parts, where a part may appear at most once. We call them “Odd Partitions of even length with Distinct parts” (ODP_{ev}); later on we also need “Odd Partitions of odd length with Distinct parts” (ODP_{odd}), their union $\text{ODP} = \text{ODP}_{\text{ev}} \cup \text{ODP}_{\text{odd}}$, and the partition 0. Hence for $0 \neq \alpha \in \text{ODP}_{\text{ev}}$ one has

$$\alpha = (\alpha_1, \alpha_2, \dots, \alpha_{2k}),$$

where all $\alpha_i \in 1 + 2\mathbb{Z}_{\geq 0}$, and we assume

$$\alpha_1 > \alpha_2 > \dots > \alpha_{2k}.$$

Introduce for such partition $0 \neq \alpha \in \text{ODP}_{\text{ev}}$

$$\xi_\alpha = t_{\frac{\alpha_1}{2}} t_{\frac{\alpha_2}{2}} \dots t_{\frac{\alpha_{2k}}{2}}, \quad \text{and} \quad \xi_0 = 1.$$

Then we can rewrite $g|0\rangle$ as

$$\sigma(g|0\rangle) = \tau(\mathbf{t}) = \sum_{\alpha \in \text{ODP}_{\text{ev}}} \tau_\alpha(t) \xi_\alpha.$$

Note that $\alpha \in \text{ODP}_{\text{ev}}$ otherwise $\langle 0|g|0\rangle = 0$. We also rewrite (2.23) as

$$\text{Res}_z \sum_{\alpha, \beta \in \text{ODP}_{\text{ev}}} \sigma\phi(z) \sigma^{-1} \tau_\alpha(t) \xi_\alpha \otimes \sigma\phi(-z) \sigma^{-1} \tau_\beta(t) \xi_\beta = 0.$$

It is clear that we can also write

$$\sigma(\phi(z_0)\phi(z_1)\dots\phi(z_k)|0\rangle) = f(\mathbf{t}) = \sum_{\alpha \in \text{ODP}, \ell(\alpha) \leq k+1} f_\alpha(t) \xi_\alpha. \quad (2.24)$$

Now let

$$\alpha = (\alpha_1, \alpha_2, \dots, \alpha_n), \quad \alpha_i > \alpha_{i+1}.$$

Then clearly

$$f_\alpha(t) = \frac{\partial}{\partial t_{\frac{\alpha_n}{2}}} \frac{\partial}{\partial t_{\frac{\alpha_{n-1}}{2}}} \dots \frac{\partial}{\partial t_{\frac{\alpha_1}{2}}} \sigma(\phi(z_0)\phi(z_1)\dots\phi(z_k)|0\rangle) \Big|_{t_{\text{odd}}=0}.$$

Now substitute the vertex operator expression for the fields $\sigma\phi(z_j)\sigma^{-1}$. We thus obtain (assuming $|z_i| > |z_j|$ if $i < j$):

$$\sigma(\phi(z_0)\phi(z_1)\dots\phi(z_k)|0\rangle) = \left(\prod_{0 \leq i < j \leq k} \frac{z_i + z_j}{z_i - z_j} \right) e^{\sum_{m=0}^k \varphi(t, z_m)} \sigma(\theta(z_0)\dots\theta(z_k)|0\rangle),$$

where

$$\varphi(t, z) = \sum_{0 < k \in \mathbb{Z}, \text{ odd}} t_k z^k. \quad (2.25)$$

Then setting $\Pi(z) := \prod_{-1 \leq i < j \leq k} \frac{z_i + z_j}{z_i - z_j}$ we have

$$f_\alpha(t) = \Pi(z) e^{\sum_{m=0}^k \varphi(t, z_m)} \frac{\partial}{\partial t_{\frac{\alpha_n}{2}}} \frac{\partial}{\partial t_{\frac{\alpha_{n-1}}{2}}} \dots \frac{\partial}{\partial t_{\frac{\alpha_1}{2}}} \sigma(\theta(z_0)\dots\theta(z_k)|0\rangle) \Big|_{t_{\text{odd}}=0}$$

$$\begin{aligned}
&= \Pi(z) e^{\sum_{m=0}^k \varphi(t, z_m)} \sigma \left(J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} \theta(z_0) \cdots \theta(z_k) |0\rangle \right) \Big|_{t_{\text{odd}}=0} \\
&= \Pi(z) e^{\sum_{m=0}^k \varphi(t, z_m)} \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} \theta(z_0) \cdots \theta(z_k) |0\rangle \\
&= \frac{\Pi(z)}{2^n} e^{\sum_{m=0}^k \varphi(t, z_m)} \text{Res}_{z_{-n}} \cdots \text{Res}_{z_{-1}} z_{-n}^{\frac{\alpha_n}{2}-\frac{1}{2}} \cdots z_{-1}^{\frac{\alpha_1}{2}-\frac{1}{2}} \langle 0 | \theta(z_{-n}) \theta(z_{-n+1}) \cdots \theta(z_k) |0\rangle \\
&= \frac{\Pi(z)}{2^n} e^{\sum_{m=0}^k \varphi(t, z_m)} \text{Res}_{z_{-n}} \cdots \text{Res}_{z_{-1}} z_{-n}^{\frac{\alpha_n}{2}-\frac{1}{2}} \cdots z_{-1}^{\frac{\alpha_1}{2}-\frac{1}{2}} \text{Pf} \left(\left(\frac{z_i - z_j}{(z_i + z_j)^2} \right)_{-n \leq i, j \leq k} \right),
\end{aligned} \tag{2.26}$$

where Pf stands for the Pfaffian. The last equality follows from Wick's theorem and from

$$\langle 0 | \theta(y) \theta(z) |0\rangle = \frac{y - z}{(y + z)^2}.$$

Define

$$\Gamma(\mathbf{t}) := e^{J(\mathbf{t})}, \quad J(\mathbf{t}) = \sum_{0 < i \in 1+2\mathbb{Z}} t_i J_i + \sum_{0 < i \in \frac{1}{2}+\mathbb{Z}} t_i J_i, \tag{2.27}$$

where $\deg t_i = i$, $\deg J(\mathbf{t}) = 0$. We write

$$\Gamma(\mathbf{t}) = e^{H(t)} e^{\chi(t_{\text{odd}})}, \quad \text{where } H(t) := \sum_{0 < i \in 1+2\mathbb{Z}} t_i J_i, \quad \chi(t_{\text{odd}}) := \sum_{0 < i \in \frac{1}{2}+\mathbb{Z}} t_i J_i. \tag{2.28}$$

We will now show that $f(\mathbf{t})$ is equal to

$$\langle 0 | \Gamma(\mathbf{t}) \phi(z_0) \phi(z_1) \cdots \phi(z_k) |0\rangle.$$

Using $e^{H(t)} \phi(z) e^{-H(t)} = e^{\varphi(t, z)}$ and $e^{\chi(t_{\text{odd}})} \phi(z) e^{-\chi(t_{\text{odd}})} = \phi(z) + \Xi(z)$, where

$$\Xi(z) = V_-(z) \xi(z) V_+(z), \quad \xi(z) = \sum_{0 < k \in \frac{1}{2}+\mathbb{Z}} k t_k (-z)^{k-\frac{1}{2}},$$

we see that

$$\begin{aligned}
&\frac{\partial}{\partial t_{\frac{\alpha_n}{2}}} \frac{\partial}{\partial t_{\frac{\alpha_{n-1}}{2}}} \cdots \frac{\partial}{\partial t_{\frac{\alpha_1}{2}}} \langle 0 | e^{H(t)} e^{\chi(t_{\text{odd}})} \phi(z_0) \phi(z_1) \cdots \phi(z_k) |0\rangle \Big|_{t_{\text{odd}}=0} \\
&= \langle 0 | e^{H(t)} J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} e^{\chi(t_{\text{odd}})} \phi(z_0) \phi(z_1) \cdots \phi(z_k) |0\rangle \Big|_{t_{\text{odd}}=0} \\
&= \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} e^{H(t)} e^{\chi(t_{\text{odd}})} \phi(z_0) \phi(z_1) \cdots \phi(z_k) |0\rangle \Big|_{t_{\text{odd}}=0} \\
&= e^{\sum_{m=0}^k \varphi(t, z_m)} \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} (\phi(z_0) + \Xi(z_0)) \\
&\quad \times (\phi(z_1) + \Xi(z_1)) \cdots (\phi(z_k) + \Xi(z_k)) |0\rangle \Big|_{t_{\text{odd}}=0} \\
&= e^{\sum_{m=0}^k \varphi(t, z_m)} \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} \phi(z_0) \phi(z_1) \cdots \phi(z_k) |0\rangle \\
&= e^{\sum_{m=0}^k \varphi(t, z_m)} \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} \theta(z_0) \theta(z_1) \cdots \theta(z_k) \\
&\quad \times V_-(z_0) V_+(z_0) V_-(z_1) V_+(z_1) \cdots V_-(z_k) V_+(z_k) |0\rangle
\end{aligned}$$

$$= \left(\prod_{0 \leq i < j \leq k} \frac{z_i + z_j}{z_i - z_j} \right) e^{\sum_{m=0}^k \varphi(t, z_m)} \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}} \theta(z_0) \theta(z_1) \cdots \theta(z_k) | 0 \rangle = f_\alpha(t).$$

Note that from this for k odd, due to Wick's rule for bosons we also have

$$\begin{aligned} f_0(t) &= e^{\sum_{m=0}^k \varphi(t, z_m)} \langle 0 | \phi(z_0) \phi(z_1) \cdots \phi(z_k) | 0 \rangle = e^{\sum_{m=0}^k \varphi(t, z_m)} \text{Hf} \left((\langle \phi(z_i) \phi(z_j) \rangle)_{0 \leq i, j \leq k} \right) \\ &= e^{\sum_{m=0}^k \varphi(t, z_m)} \text{Hf} \left(\left(\frac{1}{z_i + z_j} \right)_{0 \leq i, j \leq k} \right), \end{aligned}$$

where Hf stands for the Hafnian. The Hafnian of a symmetric matrix A of even order is defined as follows

$$\text{Hf}(A) := \sum_{\sigma} A_{\sigma(1), \sigma(2)} A_{\sigma(3), \sigma(4)} \cdots A_{\sigma(2k-1), \sigma(2k)},$$

where the sum runs over all permutations σ of $\{1, \dots, 2k\}$ satisfying

$$\sigma(2i-1) < \sigma(2i), \quad \sigma(1) < \sigma(3) < \cdots < \sigma(2k-1).$$

As one can see the Hafnian contains $1 \cdot 3 \cdot 5 \cdots (2k-1) =: (2k-1)!!$ terms.

Remark 1. Comparing this with f_0 in (2.26), we have a new proof for the identity

$$\text{Pf} \left(\left(\frac{z_i - z_j}{(z_i + z_j)^2} \right)_{1 \leq i, j \leq 2k} \right) = \prod_{1 \leq i < j \leq 2k} \frac{z_i - z_j}{z_i + z_j} \text{Hf} \left(\left(\frac{1}{z_i + z_j} \right)_{1 \leq i, j \leq 2k} \right) \quad (2.29)$$

of [7].

Since $g|0\rangle$, where g is given by (2.22), is a possibly infinite linear combination of

$$\phi_{j_1} \phi_{j_2} \cdots \phi_{j_{2k}} |0\rangle,$$

which can be obtained by taking residues of the expression in (2.24), one deduces that

$$\sigma(g|0\rangle) = \tau(\mathbf{t}) = \langle 0 | e^{H(t)} e^{\chi(t_{\text{odd}})} g | 0 \rangle = \sum_{\alpha} \tau_{\alpha}(t) \xi_{\alpha} \quad (2.30)$$

and

$$\tau_{\alpha}(t) = \langle \alpha | e^{H(t)} g | 0 \rangle, \quad (2.31)$$

where

$$\langle \alpha | = \langle 0 | J_{\frac{\alpha_n}{2}} J_{\frac{\alpha_{n-1}}{2}} \cdots J_{\frac{\alpha_1}{2}}. \quad (2.32)$$

2.6 A CKP wave function

Now we want to study $\phi(z)g|0\rangle$. Consider the expression

$$\begin{aligned} \sigma(\phi(z)g|0\rangle) &= \sigma\phi(z)\sigma^{-1}\sigma(g|0\rangle) = \sigma\phi(z)\sigma^{-1}\tau(\mathbf{t}) \\ &= e^{\varphi(t, z)} \sum_{0 < j \in \mathbb{Z}} \left((2j-1)t_{\frac{2j-1}{2}}(-z)^{j-1} + 2\frac{\partial}{\partial t_{\frac{2j-1}{2}}} z^{-j} \right) \times \end{aligned}$$

$$\times \exp \left(\sum_{0 < k \in \mathbb{Z}, \text{odd}} \frac{2}{k} \frac{\partial}{\partial t_k} z^{-k} \right) \tau(\mathbf{t}),$$

where (2.20) was used.

Clearly, one also has

$$\sigma(\phi(z)g|0\rangle) = \langle 0|e^{H(t)}e^{\chi(t_{\text{odd}})}\phi(z)g|0\rangle. \quad (2.33)$$

We now write $\sigma(\phi(z)g|0\rangle)$ as

$$\sigma(\phi(z)g|0\rangle) = \sum_{\alpha \in \text{ODP}_{\text{odd}}} g_{\alpha}(t, z) \xi_{\alpha}.$$

Now substitute this in (2.23), omitting the tensor symbol and writing s_j for t_j in the right-hand side of the tensor product, we obtain, that for every $\alpha, \beta \in \text{ODP}_{\text{odd}}$ the coefficient of $\xi_{\alpha}\eta_{\beta}$ is equal to

$$\text{Res}_z g_{\alpha}(t, z)g_{\beta}(s, -z) = 0. \quad (2.34)$$

Now we want to express $g_{\alpha}(t, z)$ in terms of the $\tau_{\beta}(t)$'s.

It will be convenient to introduce some more notation here. Let

$$\alpha = (\alpha_1, \alpha_2, \dots, \alpha_k) \in \text{ODP},$$

where all $\alpha_i \in 1 + 2\mathbb{Z}_{\geq 0}$, and we assume

$$\alpha_1 > \alpha_2 > \dots > \alpha_k.$$

Now let $\nu \in 1 + 2\mathbb{Z}_{\geq 0}$, $\nu \notin \alpha$, i.e.,

$$\nu \notin \{\alpha_1, \alpha_2, \dots, \alpha_k\},$$

and

$$\alpha_i > \nu > \alpha_{i+1},$$

then we define an “addition” as follows

$$\alpha \cup \nu := (\alpha_1, \alpha_2, \dots, \alpha_i, \nu, \alpha_{i+1}, \dots, \alpha_k).$$

Note that the notation $\alpha + \beta$ was used differently in [11] where it was defined as $(\alpha_1 + \beta_1, \alpha_2 + \beta_2, \dots)$.

In a similar way the subtraction $\alpha \setminus \alpha_i$ for $\alpha_i \in \alpha$ is defined by

$$\alpha \setminus \alpha_i = (\alpha_1, \alpha_2, \dots, \alpha_{i-1}, \alpha_{i+1}, \dots, \alpha_k).$$

Then for

$$\alpha = (\alpha_1, \alpha_2, \dots, \alpha_k) \in \text{ODP}_{\text{odd}},$$

hence k odd, we find

$$g_{\alpha}(t, z) = e^{\varphi(t, z)} \exp \left(\sum_{0 < j \in \mathbb{Z}, \text{odd}} \frac{2}{j} \frac{\partial}{\partial t_j} z^{-j} \right) \times$$

$$\times \left(\sum_{i=1}^k (-)^{i-1} \alpha_i \tau_{\alpha \setminus \alpha_i}(t) (-z)^{\frac{\alpha_i-1}{2}} + 2 \sum_{\nu \in 1+2\mathbb{Z}_{\geq 0}, \nu \notin \alpha} s(\nu, \alpha) \tau_{\alpha \cup \nu}(t) z^{-\frac{\nu+1}{2}} \right), \quad (2.35)$$

where

$$s(\nu, \alpha) = (-1)^{|\{\alpha_i \in \alpha \mid \alpha_i > \nu\}|}.$$

In particular for $\alpha = (1) = 1$ we find

$$g_1(t, z) = e^{\varphi(t, z)} \exp \left(\sum_{0 < j \in \mathbb{Z}, \text{odd}} \frac{2}{j} \frac{\partial}{\partial t_j} z^{-j} \right) \left(\tau_0(t) + 2 \sum_{\nu \in 1+2\mathbb{Z}, \nu > 1} \tau_{(\nu, 1)}(t) z^{-\frac{\nu+1}{2}} \right),$$

We now want to calculate $g_\alpha(t, z)$ as some expectation value. Using (2.33), we have

$$\begin{aligned} g_\alpha(t, z) &= \frac{\partial}{\partial t_{\frac{\alpha_n}{2}}} \frac{\partial}{\partial t_{\frac{\alpha_{n-1}}{2}}} \cdots \frac{\partial}{\partial t_{\frac{\alpha_1}{2}}} \langle 0 | e^{H(t)} e^{\chi(t_{\text{odd}})} \phi(z) g | 0 \rangle \Big|_{t_{\text{odd}}=0} \\ &= \frac{1}{2^n} \langle 0 | \theta_{\frac{\alpha_n}{2}} \theta_{\frac{\alpha_{n-1}}{2}} \cdots \theta_{\frac{\alpha_1}{2}} e^{H(t)} \phi(z) g | 0 \rangle = \langle \alpha | e^{H(t)} \phi(z) g | 0 \rangle. \end{aligned} \quad (2.36)$$

where we have used (2.32).

Now concentrating on (2.35) and divide this by $(-)^{\frac{\alpha_1-1}{2}} \alpha_1 \tau_{\alpha \setminus \alpha_1}(t)$, this gives

$$\begin{aligned} w_\alpha(t, z) &= \hat{w}_\alpha(t, z) z^{\frac{\alpha_1-1}{2}} e^{\varphi(t, z)}, \quad \text{where} \\ \hat{w}_\alpha(t, z) &= (-z)^{-\frac{\alpha_1-1}{2}} \frac{g_\alpha(t, z) e^{-\varphi(t, z)}}{\alpha_1 \tau_{\alpha \setminus \alpha_1}(t)} = 1 + O(z^{-1}), \end{aligned} \quad (2.37)$$

which we call the wave function corresponding to $\alpha \in \text{ODP}$. Now using (2.31) and (2.36), one also has

$$w_\alpha(t, z) = \frac{(-)^{\frac{\alpha_1-1}{2}}}{\alpha_1} \frac{\langle \alpha | e^{H(t)} \phi(z) g | 0 \rangle}{\langle \alpha \setminus \alpha_1 | g | 0 \rangle}.$$

2.7 A bilinear identity for the wave function and a Lax equation

Using the definition of the wave function (2.37) we can rewrite (2.34) into a bilinear identity for these wave functions:

$$\begin{aligned} \text{Res}_z w_\alpha(t, z) w_\beta(s, -z) \\ = \text{Res}_z \hat{w}_\alpha(t, z) \hat{w}_\beta(s, -z) e^{\varphi(t, z) - \varphi(s, z)} z^{\frac{\alpha_1 + \beta_1 - 2}{2}} = 0, \quad \alpha, \beta \in \text{ODP}_{\text{odd}}. \end{aligned} \quad (2.38)$$

Let us take $\alpha = \beta$ and $\alpha_1 = \beta_1 = n$, n odd. Then equation (2.38) takes the form

$$\text{Res}_z \hat{w}_\alpha(t, z) \hat{w}_\alpha(s, -z) e^{\varphi(t, z) - \varphi(s, z)} z^{n-1} = 0, \quad \alpha_1 = \beta_1 = n, \quad n \text{ odd},$$

which coincides with equations (5) and $(12)_n$ (n odd) of [3] where $w(x, k)$ should be compared with our $e^{\varphi(t, z)} w_\alpha(t, z)$.

Now write for $\alpha = (1) = 1$

$$w(t, z) = w_1(t, z), \quad \text{and} \quad \hat{w}(t, z) = \hat{w}_1(t, z)$$

then we can see this as the Date, Jimbo, Kashiwara, Miwa CKP wave function as in [3], since in particular

$$\text{Res}_z \hat{w}(t, z) \hat{w}(s, -z) e^{\varphi(t, z) - \varphi(s, z)} = 0. \quad (2.39)$$

We can rewrite the wave functions as follows

$$w_\alpha(t, z) = \hat{w}_\alpha(t, \partial) \partial^{\frac{\alpha_1-1}{2}} e^{\varphi(t, z)},$$

where $\partial = \frac{\partial}{\partial x} = \frac{\partial}{\partial t_1}$. Then $\hat{w}_\alpha(t, \partial)$ is a pseudo differential operator of order 0. Note that

$$w_\alpha(t, -z) = \hat{w}_\alpha(t, \partial) \partial^{\frac{\alpha_1-1}{2}} e^{-\varphi(t, z)}.$$

Now we use the following known lemma (see, e.g., [10]):

Lemma 1. *For pseudodifferential operators $P(t, \partial)$ and $Q(t, \partial)$ we have*

$$\begin{aligned} \text{Res}_z \left(P(t, \partial) \cdot \exp \left(\sum_{0 < j \in \mathbb{Z}, \text{odd}} t_j z^j \right) \right) \left(Q(t, \partial) \cdot \exp \left(- \sum_{0 < j \in \mathbb{Z}, \text{odd}} t_j z^j \right) \right) \\ = \text{Res}_\partial P(t, \partial) \cdot Q(t, \partial)^* = 0. \end{aligned}$$

where the conjugation and Res_∂ are defined on monomials respectively as $(a(x)\partial^k)^* = (-\partial)^k a(x)$ and $\text{Res}_\partial a(x)\partial^k = a(x)\delta_{k,-1}$.

Taking $P = \hat{w}$ and $Q = \partial^n \hat{w}$, $n = 0, 1, 2, \dots$, by (2.39) we obtain that

$$\left(\hat{w}_\alpha(t, \partial) \partial^{\frac{\alpha_1+\beta_1-2}{2}} \hat{w}_\beta(t, \partial)^* \right)_- = 0. \quad (2.40)$$

Now we take $\alpha = \beta = (1) = 1$. Then, since $\hat{w}(t, \partial) = 1 + \sum_{k=1}^{\infty} w_k(t) \partial^{-k}$, one obtains that

$$\hat{w}(t, \partial)^* = \hat{w}(t, \partial)^{-1}. \quad (2.41)$$

Start again with (2.39) and differentiate this equation in t_k for k odd; then one gets

$$\text{Res}_z \left(\frac{\partial \hat{w}(t, z)}{\partial t_k} + \hat{w}(t, z) z^k \right) \exp \left(\sum_{0 < j \in \mathbb{Z}, \text{odd}} t_j z^j \right) w(s, -z) = 0.$$

Again using the fundamental Lemma and (2.41), we deduce

$$\left(\left(\frac{\partial \hat{w}(t, \partial)}{\partial t_k} + \hat{w}(t, \partial) \partial^k \right) W(t, \partial)^{-1} \right)_- = 0,$$

which gives the Sato–Wilson equation:

$$\frac{\partial \hat{w}(t, \partial)}{\partial t_k} = - \left(\hat{w}(t, \partial) \partial^k \hat{w}(t, \partial)^{-1} \right)_- \hat{w}(t, \partial).$$

This is equivalent to the following equation for the wave function

$$\frac{\partial w(t, z)}{\partial t_k} = \left(\hat{w}(t, \partial) \partial^k \hat{w}(t, \partial)^{-1} \right)_+ w(t, z).$$

Introducing the Lax operator

$$L(t, \partial) = \hat{w}(t, \partial) \partial \hat{w}(t, \partial)^{-1} = \partial + \sum_{k=1}^{\infty} u_k(t) \partial^{-k},$$

then (2.41) induces $L^* = -L$ and from the Sato–Wilson equation one derives the following CKP Lax equation:

$$\frac{\partial L(t, \partial)}{\partial t_k} = \left[(L(t, \partial)^k)_+, L(t, \partial) \right].$$

3 Bilinear identity in super notations

In the previous section we obtained a wave function. In this section our approach will be slightly different. We want to superize, i.e., obtain a supersymmetric wave function that also include the Grassmannian times t_i with $i \in \frac{1}{2} + \mathbb{Z}$ and the corresponding bilinear equation for this super wave function (3.11). Let us point out that in this way we shall re-write results of the previous section using super notations. We regard this an important step, which might be very fundamental for the further development of the theory. However, unfortunately we were not able to obtain Lax equations with respect to odd Grassmannian times in this setting.

3.1 Super vertex operator

Recall the super commutation relations (2.18) and the super times (2.18) and the definition of $\Gamma(\mathbf{t})$ in (2.27). It is convenient to introduce an auxiliary parameter ζ , which is a Grassmannian variable, an odd counterpart to z : $\zeta^2 = 0$, $z\zeta = \zeta z$, $\deg z = 2 \deg \zeta = -1$.

Introduce the following “superfermionic” fields

$$\begin{aligned}\Phi(z, \zeta) &:= 2 \sum_{n \in \mathbb{Z}} \zeta \frac{z^{-2n-1}}{2n+1} J_{2n+1} - 2 \sum_{0 \neq n \in \mathbb{Z}} \frac{z^{-n}}{n} J_{n+\frac{1}{2}} + 2J_{\frac{1}{2}} \log z, \\ \Theta(z, \zeta) &:= \left(\frac{\partial}{\partial \zeta} + \zeta \frac{\partial}{\partial z} \right) \Phi(z, \zeta).\end{aligned}$$

Here $\frac{\partial}{\partial \zeta} + \zeta \frac{\partial}{\partial z}$ is a superderivative which will be denoted by $D_{z, \zeta}$, $D_{z, \zeta}^2 = \frac{\partial}{\partial z}$. As we see $\deg \Phi = -\frac{1}{2}$ while $\deg \Theta = 0$.

Let $:e^{\Theta(z, \zeta)}:$ denote $e^{\Theta_-(z, \zeta)} e^{\Theta_+(z, \zeta)}$ where Θ_{\pm} denotes the splitting of Θ in series in respectively positive/negative powers of z^{-1} . One can verify

$$:e^{\Theta(z_1, \zeta_1)}: :e^{\Theta(z_2, \zeta_2)}: = e^{\Theta_-(z_1, \zeta_1) + \Theta_-(z_2, \zeta_2)} e^{\Theta_+(z_1, \zeta_1) + \Theta_+(z_2, \zeta_2)} \left(\frac{z_1 + z_2}{z_1 - z_2} + \frac{\zeta_1 \zeta_2}{z_1 + z_2} \right).$$

It follows from $\zeta^2 = 0$ that

$$:e^{\Theta(z, \zeta)}: = V_-(z) V_+(z) + \zeta \phi(z). \quad (3.1)$$

Then

$$\begin{aligned}\phi(z) &= \frac{\partial}{\partial \zeta} :e^{\Theta(z, \zeta)}: \\ &= \frac{\partial}{\partial \zeta} \left(\sum_{e^{0 < n \in \mathbb{Z}}} \left(\frac{2z^{-2n-1}}{2n+1} J_{2n+1} + 2\zeta z^{-1-n} J_{n+\frac{1}{2}} \right) \cdot e^{\sum_{0 \leq n \in \mathbb{Z}} \left(\frac{2z^{-2n-1}}{2n+1} J_{2n+1} + 2\zeta z^{-n-1} J_{n+\frac{1}{2}} \right)} \right). \quad (3.2)\end{aligned}$$

Using the isomorphism σ , see (2.19), we can express

$$\sigma \Theta_-(z, \zeta) \sigma^{-1} = \sum_{0 < i \in 1+2\mathbb{Z}} z^i t_i + \zeta \sum_{0 < i \in \frac{1}{2} + \mathbb{Z}} i(-z)^{i-\frac{1}{2}} t_i =: \varphi(\mathbf{t}, z, \zeta), \quad (3.3)$$

and

$$\sigma \Theta_+(z, \zeta) \sigma^{-1} = 2 \sum_{0 < n \in \mathbb{Z}} \frac{z^{1-2n}}{2n-1} \frac{\partial}{\partial t_{2n-1}} + 2\zeta \sum_{0 < n \in \mathbb{Z}} z^{-n} \frac{\partial}{\partial t_{n-\frac{1}{2}}}.$$

Now

$$\sigma\phi(z)\sigma^{-1} = \frac{\partial}{\partial\zeta}\Upsilon(\mathbf{t}, z, \zeta), \quad (3.4)$$

where

$$\begin{aligned} \Upsilon(\mathbf{t}, z, \zeta) &:= e^{\varphi(\mathbf{t}, z, \zeta)}\Upsilon_+(z, \zeta), \\ \Upsilon_+(z, \zeta) &= \exp\left(2\sum_{0 < i \in 1+2\mathbb{Z}} \frac{z^{-i}}{i} \frac{\partial}{\partial t_i} + 2\zeta \sum_{0 < i \in \frac{1}{2}+\mathbb{Z}} z^{-i-\frac{1}{2}} \frac{\partial}{\partial t_i}\right). \end{aligned} \quad (3.5)$$

3.2 Tau function and super wave function

Recall (2.30), (2.31) the tau function as a function of a collection of both even and odd time variables as follows

$$\tau(\mathbf{t}) := \langle 0 | \Gamma(\mathbf{t}) g | 0 \rangle = \sum_{\alpha} \tau_{\alpha}(t) \xi_{\alpha}. \quad (3.6)$$

As we can see

$$\tau(\mathbf{t}) = \tau^{\text{CKP}}(t) + \dots = \tau_0(t) + \dots,$$

where the dots mean terms vanishing when we substitute $t_{\frac{1}{2}} = t_{\frac{3}{2}} = \dots = 0$. We put

$$W(\mathbf{t}, z, \zeta) := \frac{1}{\tau(\mathbf{t})} \langle 0 | \Gamma(\mathbf{t}) : e^{\Theta(z, \zeta)} : g | 0 \rangle \quad (3.7)$$

to be the wave function.

From (3.4) we can write

$$W(\mathbf{t}, z, \zeta) := \frac{\Upsilon(\mathbf{t}, z, \zeta)\tau(\mathbf{t})}{\tau(\mathbf{t})} = e^{\varphi(\mathbf{t}, z, \zeta)} \left(1 + O\left(\frac{1}{z}\right)\right), \quad (3.8)$$

where $\varphi(\mathbf{t}, z, \zeta)$ was defined in (3.3). Let us mark that if $t_{\text{odd}} = 0$ then $\varphi(\mathbf{t}, z, \zeta)$ coincides with $\varphi(t, z, \zeta)$ defined in (2.25).

Note that since $\phi(z) = \frac{\partial}{\partial\zeta} : e^{\Theta(z, \zeta)} :$ it follows from (2.35) and (2.36) that for $\alpha = (\alpha_1, \dots, \alpha_n)$

$$g_{\alpha}(t, z) = \left(\frac{\partial}{\partial\zeta} \frac{\partial}{\partial t_{\frac{\alpha_n}{2}}} \cdots \frac{\partial}{\partial t_{\frac{\alpha_1}{2}}} \Upsilon(\mathbf{t}, \zeta, z) \tau(\mathbf{t}) \right) \Big|_{t_{\text{odd}}=0},$$

where we remind that in our notations: $t = (t_1, t_3, t_5, \dots)$ and $\mathbf{t} = (t_1, t_{\frac{1}{2}}; t_3, t_{\frac{3}{2}}; t_5, t_{\frac{5}{2}}; \dots)$.

3.3 Super Miwa variables

Denote $\mathbf{z} = (z_1, \zeta_1; \dots; z_k, \zeta_k)$, where ζ_i are Grassmannian odd and z_i are Grassmannian even variables, $\zeta_i^2 = 0$, and $\deg z_i = -1$, $\deg \zeta_i = -\frac{1}{2}$. Let $[\mathbf{z}] := (t_1, t_{\frac{1}{2}}; t_3, t_{\frac{3}{2}}; t_5, t_{\frac{5}{2}}; \dots)$, where

$$t_{2n+1} = \frac{2}{(2n+1)} \sum_{i=1}^k \frac{1}{z_i^{2n+1}}, \quad t_{n+\frac{1}{2}} = 2 \sum_{i=1}^k \frac{\zeta_i}{z_i^n}. \quad (3.9)$$

For $k = 1$ we shall write $[\mathbf{z}]$ as $[z, \zeta]$. Then

$$W(\mathbf{t}, z, \zeta) := e^{\varphi(\mathbf{t}, z, \zeta)} \frac{\tau(\mathbf{t} - [z, \zeta])}{\tau(\mathbf{t})}.$$

In general

$$\langle 0 | \Gamma(\mathbf{t} - [\mathbf{z}]) = \mathbb{D}(\mathbf{z})^{-1} \langle 0 | \prod_{i=1}^k :e^{\Theta(z_i, \zeta_i)}: \Gamma(\mathbf{t}),$$

where

$$\mathbb{D}(\mathbf{z}) = \mathbb{D}_k(\mathbf{z}) := \langle 0 | :e^{\Theta(z_1, \zeta_1)}: \dots :e^{\Theta(z_k, \zeta_k)}: | 0 \rangle,$$

which is antisymmetric function with respect to the permutation of pairs (z_i, ζ_i) . For instance

$$\mathbb{D}_2(\mathbf{z}) = \frac{z_1 + z_2}{z_1 - z_2} + \frac{\zeta_1 \zeta_2}{z_1 + z_2}.$$

We have

$$\mathbb{D}_k(\mathbf{z})^{-1} = \left(\prod_{a < b} \frac{z_a - z_b}{z_a + z_b} \right) \sum_{n=0}^{[k/2]} (-1)^n \sum_{\substack{\alpha \in \text{ODP}_{\text{ev}} \\ \ell(\alpha) = 2n}} \zeta_\alpha \text{Pf} \left[\frac{z_{\alpha_i} - z_{\alpha_j}}{(z_{\alpha_i} + z_{\alpha_j})^2} \right]. \quad (3.10)$$

3.4 Bilinear identity for the super wave function W

The bilinear equation (2.34),

$$\text{Res}_z g_\alpha(t', z) g_\beta(t, -z) = 0,$$

may be written in superized form

$$\text{Ber}_{z, \zeta} \left(\left(\frac{\partial}{\partial \zeta} W(\mathbf{t}', z, \zeta) \right) W(\mathbf{t}, -z, \mp \zeta) \right) = 0, \quad (3.11)$$

where $\text{Ber}_{z, \zeta} f(z, \zeta) := \text{Res}_z \frac{\partial f(z, \zeta)}{\partial \zeta}$. The validity of (3.11) follows directly from (2.23), (3.7), (3.2) and the definition of the Berezinian Ber .

3.5 Bilinear identity as identities for super PDOs

Let us re-write super-bilinear identities (3.11) in form of identities for (super version of) pseudo-differential operators (PDO). We will do it in a way similar to the KP case where PDO operators naturally appear in the framework of KP bilinear identity [4].

Now notice that (3.8) yields

$$W(\mathbf{t}, z, \zeta) =: \hat{W}(\mathbf{t}, z, \zeta) e^{\varphi(\mathbf{t}, z, \zeta)} =: \left(\hat{W}^{(0)}(\mathbf{t}, z) + \zeta \hat{W}^{(1)}(\mathbf{t}, z) \right) e^{\varphi(\mathbf{t}, z, \zeta)}, \quad (3.12)$$

where (3.5) provides

$$\hat{W}^{(1)}(\mathbf{t}, z) = \sum_{n > 0} \hat{W}_n^{(1)}(\mathbf{t}) z^{-n} = O\left(\frac{1}{z}\right), \quad \hat{W}^{(0)}(\mathbf{t}, z) = \sum_{n \geq 0} \hat{W}_n^{(0)}(\mathbf{t}) z^{-n} = 1 + O\left(\frac{1}{z}\right).$$

For a simplification we shall denote $x = t_1$ and $\xi = \frac{1}{2}t_{\frac{1}{2}}$, thus $\varphi(\mathbf{t}, z, \zeta) = zx + \zeta\xi + \dots$, and we have

$$\partial_x e^{\varphi(\mathbf{t}, z, \zeta)} = z e^{\varphi(\mathbf{t}, z, \zeta)}, \quad \partial_\xi e^{\varphi(\mathbf{t}, z, \zeta)} = -\zeta e^{\varphi(\mathbf{t}, z, \zeta)}.$$

Let $D := D_{x, \xi} = \partial_\xi + \xi \partial_x$, $D^2 = \partial_x$. Then $\partial_\xi = D - \xi D^2$. Below we consider the action of odd negative powers of D on the exponentials $e^{zx + \zeta\xi}$, provided we define $\partial^n e^{zx} := z^n e^{zx}$, $n \in \mathbb{Z}$. To do this we write $D^{1-2n} = D \partial^{-n}$. In such a way we write

$$D^{2n} \cdot e^{\varphi(\mathbf{t}, z, \zeta)} = z^n e^{\varphi(\mathbf{t}, z, \zeta)}, \quad D^{2n+1} \cdot e^{\varphi(\mathbf{t}, z, \zeta)} = z^n (-\zeta + \xi D^2) e^{\varphi(\mathbf{t}, z, \zeta)}, \quad n \in \mathbb{Z}; \quad (3.13)$$

in particular

$$\zeta e^{\varphi(\mathbf{t}, z, \zeta)} = -(D - \xi D^2) \cdot e^{\varphi(\mathbf{t}, z, \zeta)}. \quad (3.14)$$

Let us introduce

$$\begin{aligned} K(\mathbf{t}, D) &:= \sum_{n \geq 0} K_n(\mathbf{t}) D^{-n} := \hat{W}^{(0)}(\mathbf{t}, \partial) - \hat{W}^{(1)}(\mathbf{t}, \partial) \frac{\partial}{\partial \xi} \\ &= \hat{W}^{(0)}(\mathbf{t}, D^2) - \hat{W}^{(1)}(\mathbf{t}, D^2) (D - \xi D^2) \\ &= 1 + \hat{W}_1^{(1)}(\mathbf{t}) \xi + \sum_{n \geq 1} \left(\hat{W}_n^{(0)}(\mathbf{t}) + \hat{W}_{n+1}^{(1)}(\mathbf{t}) \xi \right) D^{-2n} - \sum_{n \geq 1} \hat{W}_n^{(1)}(\mathbf{t}) D^{1-2n}, \end{aligned} \quad (3.15)$$

where negative powers of D are to be understood in the sense of (3.13).

From (3.14), (3.15) and (3.12) it follows that

$$W(\mathbf{t}, z, \zeta) = K(\mathbf{t}, D) \cdot e^{\varphi(\mathbf{t}, z, \zeta)}.$$

In the Lemma below star means the conjugation in the algebra of super PDOs with properties $(ab)^* = \pm b^* a^*$, where $-$ is taken iff both a and b are odd. We have $(\partial_x)^* = -\partial_x$, $(\partial_\xi)^* = -\partial_\xi$. For instance $D^* = -\partial_\xi + (-\partial_x)\xi = -D$, $(D^2)^* = -D^* D^* = -\partial$, $(\xi\zeta)^* = -\zeta\xi = \xi\zeta$, and $(aD^n)^* = (-1)^{[n/2]} D^n a$ if a is odd.

We define $\text{Res}_D \sum_{n \in \mathbb{Z}} f_n(x, \xi) D^n := f_{-1}(x, \xi)$. Notice that

$$\text{Res}_D f(x, \xi) D^{-1} g(x, \xi) = \pm \text{Res}_D f(x, \xi) g(x, \xi) D^{-1}, \quad (3.16)$$

where $+$ and $-$ are taken if g is respectively even and odd.

Lemma 2. *Let $P(D) = \sum_n P_n(x, \xi) D^n$ be an odd and $Q(D) = \sum_n Q_n(x, \xi) D^n$ an even PDO. Then*

$$\text{Ber}_{z, \zeta} (P(D) \cdot e^{zx + \zeta\xi}) (Q(D) \cdot e^{-zx - \zeta\xi}) = \text{Res}_D P(D) (Q(D))^*, \quad (3.17)$$

$$\text{Ber}_{z, \zeta} (P(D) \cdot e^{zx + \zeta\xi}) (Q(D) \cdot e^{-zx + \zeta\xi}) = \text{Res}_D P(D) (Q(-D))^*. \quad (3.18)$$

Proof. Let

$$P(D) = \sum_{n \in \mathbb{Z}} \left(P_n^{(0)} D^{2n} + P_n^{(1)} D^{2n} D^{-1} \right), \quad Q(D) = \sum_{n \in \mathbb{Z}} \left(Q_n^{(0)} D^{2n} + Q_n^{(1)} D^{2n} D^{-1} \right),$$

where $P_n^{(0)}, Q_n^{(1)}$ are odd and $P_n^{(1)}, Q_n^{(0)}$ are even. Then

$$(Q(D))^* = \sum_{n \in \mathbb{Z}} \left((-1)^n D^{2n} Q_n^{(0)} - (-1)^n D^{-1} D^{2n} Q_n^{(1)} \right),$$

where we used $(D^*)^{2n} = (-1)^n D^{2n}$, $(D^{-1})^* = (D^{-2} D)^* = (D)^{-1}$.

Using (3.16), we find that the right-hand side of (3.17) is equal to

$$\sum_{n \in \mathbb{Z}} (-1)^n \left(P_n^{(1)} Q_{-n}^{(0)} + P_n^{(0)} Q_{-n}^{(1)} \right). \quad (3.19)$$

Next consider the left-hand side of (3.17). We have

$$\begin{aligned} P(D) \cdot e^{zx+\zeta\xi} &= \sum_{n \in \mathbb{Z}} \left(P_n^{(0)} z^n + P_n^{(1)} z^{n-1} (\zeta + z\xi) \right) e^{zx+\zeta\xi}, \\ Q(D) \cdot e^{-zx-\zeta\xi} &= \sum_{n \in \mathbb{Z}} \left(Q_n^{(0)} (-z)^n - Q_n^{(1)} (-z)^{n-1} (\zeta + z\xi) \right) e^{-zx-\zeta\xi}. \end{aligned}$$

The evaluation of the Ber of the product of these two results in (3.19).

A similar calculation yields (3.18). ■

We re-write (3.11) as

$$\text{Ber}_{z,\zeta} \left(K(\mathbf{t}', D) \cdot \frac{\partial}{\partial \zeta} e^{\varphi(\mathbf{t}', z, \zeta)} \right) \left(K(\mathbf{t}, D) \cdot e^{\varphi(\mathbf{t}, -z, \mp \zeta)} \right) = 0.$$

Taking into account

$$\frac{\partial}{\partial \zeta} e^{\varphi(\mathbf{t}, z, \zeta)} = - \sum_{n \geq 0} \left((-1)^n \left(n + \frac{1}{2} \right) t_{n+\frac{1}{2}} \partial^n \right) \cdot e^{\varphi(\mathbf{t}, z, \zeta)},$$

where $\partial := \frac{\partial}{\partial t_1} = D^2$, and $\varphi(\mathbf{t}, z, \zeta)$ is as in (3.3), we obtain

$$\begin{aligned} \text{Res}_D \left(K(\mathbf{t}', D) \cdot \left(\sum_{n \geq 0} (-1)^n \left(n + \frac{1}{2} \right) t'_{n+\frac{1}{2}} D^{2n} \right) \right. \\ \left. \times e^{\pm(D-\xi D^2) \sum_{n \geq 0} (n+\frac{1}{2}) t_{n+\frac{1}{2}} D^{2n}} \cdot K^*(\mathbf{t}, \pm D) \right) = 0, \end{aligned}$$

which results in

$$\left(K(\mathbf{t}, D) \cdot \left(\sum_{n \geq 0} (-1)^n \left(n + \frac{1}{2} \right) t_{n+\frac{1}{2}} D^{2n} \right) e^{\pm(D-\xi D^2) \sum_{n \geq 0} (n+\frac{1}{2}) t_{n+\frac{1}{2}} D^{2n}} \cdot K^*(\mathbf{t}, \pm D) \right)_{-} = 0,$$

where the subscript $-$ means the taking of projection on series with negative powers. Thus

$$L_{\pm} := K(\mathbf{t}, D) \cdot \left(\sum_{n \geq 0} (-1)^n \left(n + \frac{1}{2} \right) t_{n+\frac{1}{2}} D^{2n} \right) e^{\pm(D-\xi D^2) \sum_{n \geq 0} (n+\frac{1}{2}) t_{n+\frac{1}{2}} D^{2n}} \cdot K^*(\mathbf{t}, \pm D)$$

are differential operators. These equations are basically equivalent to the set of equations (2.40).

4 Related symmetric functions

In this section we want to introduce polynomial functions, C_{λ} , related to the basis vectors $|\lambda\rangle$ of the bosonic Fock space F as the image of the mapping σ described in the Subsection 2.3. These are polynomials in Grassmannian even and odd variables \mathbf{t} . In super Miwa variables $z_i, \zeta_i, i = 1, \dots, k$ these functions are symmetric with respect to the action of the permutation group S_k on the set of pairs (z_i, ζ_i) , and polynomial with respect to variables ζ_i and $x_i := z_i^{-1}$.

These polynomials may be considered as analogues of the celebrated Schur (and projective Schur) functions which are related to Fock space of charged (resp. neutral) fermions. The theoretic field construction of new functions allows to derive certain properties which are similar to the properties of the Schur and the projective Schur functions.

4.1 Polynomials \mathbf{C}_λ

Let us introduce suitable notations for the basis of bosonic Fock vectors in F (see (2.3)) and in F^* labeled by partitions whose parts are odd numbers

$$|\lambda\rangle := \frac{1}{d_\lambda} \phi_{\frac{\lambda_1}{2}} \cdots \phi_{\frac{\lambda_k}{2}} |0\rangle, \quad \langle\lambda| := \frac{1}{d_\lambda} \langle 0| \phi_{-\frac{\lambda_k}{2}} \cdots \phi_{-\frac{\lambda_1}{2}}, \quad (4.1)$$

where $\lambda = (\lambda_1, \lambda_2, \dots, \lambda_k)$ is a set of odd numbers and $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_k > 0$, $\ell(\lambda) := k = 1, 2, \dots$. The set of partitions with odd parts will be denoted by OP.

Note that the above vectors $|\lambda\rangle$ differ from the vector $\langle\alpha|$ as defined in (2.32), that is the reason why we write λ here to avoid this confusion.

In this section we shall use, besides the parts λ_i of partitions λ , also the variables $n_i = 0, 1, 2, \dots$ related to odd numbers λ_i as follows

$$\lambda_i =: 2n_i + 1. \quad (4.2)$$

The sum $|\lambda| := \lambda_1 + \dots + \lambda_k$ is called the weight of the partition λ . Here the factor d_λ for $\lambda \in \text{OP}$ is defined by

$$(d_\lambda)^2 := (-1)^{\sum_{i=1}^k n_i} \prod_{i=1,2,3,\dots} m_{2i-1}! = \prod_{i=1,2,3,\dots} m_{2i-1}! (-1)^{m_{4i-1}} = (-1)^{\frac{1}{2}(|\lambda| - \ell(\lambda))} \prod_{i=1,3,5,\dots} m_i!,$$

where $m_i = m_i(\lambda)$ is the number of parts of λ equal to i (or, the same, the multiplicity of i). (Then we can denote the partition by its frequency notation $\lambda = (1^{m_1} 3^{m_3} 5^{m_5} \dots)$). For instance $d_{(1^3 3^2 5^1)} = 3! 2! (-1)^{m_3} 1!$.

From (2.1),

$$\phi_{-\frac{1}{2}\lambda_i} \phi_{\frac{1}{2}\lambda_j} - \phi_{\frac{1}{2}\lambda_j} \phi_{-\frac{1}{2}\lambda_i} = (-1)^{n_j} \delta_{\lambda_i, \lambda_j},$$

we see that, for partitions λ and μ , we have the ortho-normality condition

$$\langle\lambda|\mu\rangle = \delta_{\mu,\lambda}, \quad \lambda, \mu \in \text{OP}. \quad (4.3)$$

Let us introduce the following functions

$$\mathbf{C}_\lambda(\mathbf{t}) := \sigma(|\lambda\rangle) := \langle 0|\Gamma(\mathbf{t})|\lambda\rangle, \quad \lambda \in \text{OP}. \quad (4.4)$$

These functions are weighted polynomials of weight $\frac{1}{2}|\lambda|$ in the variables $\mathbf{t} := (t_1, t_{\frac{1}{2}}, t_3, t_{\frac{3}{2}}, \dots)$ where $\deg t_j := j$. For example

$$\mathbf{C}_{(1)} = t_{\frac{1}{2}}, \quad \mathbf{C}_{(1^2)} = \frac{t_1}{\sqrt{2}}.$$

We evaluate \mathbf{C}_λ in case all (Grassmannian) odd variables vanish: $t_{n+\frac{1}{2}} = 0$, $n = 0, 1, 2, \dots$. (Recall that in this case we denote \mathbf{t} as $t = (t_1, t_3, \dots)$.)

Remark 2. Weighted polynomial functions are often presented as symmetric functions of some variables x_1, \dots, x_N , where the number of variables may be irrelevant. In our case we put

$$t_n = \frac{1}{n} \sum_{i=1}^N (x_i^n - (-x_i)^n), \quad n = 1, 2, 3, \dots, \quad (4.5)$$

where all even-labeled t_n vanish. For n odd we write

$$t_n = \frac{2}{n} \sum_{i=1}^N x_i^n, \quad n = 1, 3, \dots,$$

where $x_i = \frac{1}{z_i}$, $i = 1, \dots, N$. Below by polynomial functions in Miwa variables we mean polynomial functions in the variables $x_i = \frac{1}{z_i}$.

Then $C_\lambda(t)$ vanishes if $\ell(\lambda)$ is odd¹. If $\ell(\lambda)$ is even, then by Wick's theorem

$$C_\lambda(t) = \text{Hf} [C_{(\lambda_i, \lambda_j)}(t)], \quad \lambda \in \text{OP}_e,$$

where OP_e is the set of all partitions with even number of odd parts, and as we shall see

$$C_{(\lambda_i, \lambda_j)}(t) = \frac{1}{d_{(\lambda_i | \lambda_j)}} s_{(n_i | n_j)}(t), \quad (4.6)$$

where s_λ is the Schur function, and $(n | m)$ is a one-hook partition in the Frobenius notations, see [11]. Indeed, first, from

$$[J_{2n-1}, \phi_i] = \phi_{i-2n+1}$$

(cf. (2.9)) it follows

$$\phi_i(t) := \Gamma(t) \phi_i \Gamma(t)^{-1} = \sum_{n=0}^{\infty} h_n(t) \phi_{i-n},$$

where h_n are complete symmetric functions [11]. Now from (2.1), (2.2) we obtain

$$\langle 0 | \phi_{n_1 + \frac{1}{2}}(t) \phi_{n_2 + \frac{1}{2}}(t) | 0 \rangle = \sum_{n=0}^{\lambda_1} (-1)^{n_2 - n} h_n(t) h_{n_1 + n_2 + 1 - n}(t),$$

while Schur function evaluated on a one-hook partition is (see Chapter I, § 3, Example 9 in [11])

$$s_{(n_1 | n_2)}(t) = h_{n_1+1}(t) e_{n_2}(t) - h_{n_1+2}(t) e_{n_2-1}(t) + \dots + (-1)^{n_2} h_{n_1+n_2+1}(t),$$

where e_n are elementary symmetric functions. Then taking into account that for $t = (t_1, t_3, \dots)$ of form (4.5) we get the equality $h_m(t) \equiv e_m(t)$, we obtain (4.6).

Thus we get

$$C_\lambda(t) = \begin{cases} \frac{1}{d_\lambda} \text{Hf} [s_{(n_i | n_j)}(t)] & \text{if } \ell(\lambda) \text{ even,} \\ 0 & \text{if } \ell(\lambda) \text{ odd.} \end{cases} \quad (4.7)$$

It follows from (4.7) and from $s_\lambda(-t) = (-1)^{|\lambda|} s_{\lambda^{tr}}(t)$ that

$$C_\lambda(-t) = (-1)^{\frac{1}{2}(|\lambda| + \ell(\lambda))} C_\lambda(t). \quad (4.8)$$

Next, if all Grassmannian odd variables except $t_{\frac{1}{2}}$ vanish, we obtain

$$C_\lambda(t) = \begin{cases} \frac{1}{d_\lambda} \text{Hf} [s_{(n_i | n_j)}(t)] & \text{if } \ell(\lambda) \text{ even,} \\ t_{\frac{1}{2}} \frac{1}{d_\lambda} \text{Hf} [\tilde{S}] & \text{if } \ell(\lambda) = 2n - 1 \text{ odd,} \end{cases}$$

¹We need to make difference between odd numbers and odd numbers in the Grassmannian sense. In the last case we shall necessarily say Grassmannian odd number.

where \tilde{S} is $2n \times 2n$ symmetric matrix

$$\tilde{S}_{ij} = \tilde{S}_{ji} := \begin{cases} s_{(n_i|n_j)}(t) & \text{if } 1 \leq i < j \leq 2n-1, \\ s_{(n_i)}(t) & \text{if } 1 \leq i < j = 2n. \end{cases}$$

Recall that $t = (t_1, t_3, t_5, \dots)$ and that λ_i are related to n_i via (4.2).

In general case, where the odd Grassmannian variables do not vanish, we can see that C_λ is of even Grassmannian parity in case $\ell(\lambda)$ is even, and it is of odd Grassmannian parity in case $\ell(\lambda)$ is odd. We can write

$$C_\lambda(t) = \begin{cases} \frac{1}{d_\lambda} \text{Hf} [s_{(n_i|n_j)}(t)] + C_e & \text{if } \ell(\lambda) \text{ even,} \\ t_{\frac{1}{2}} \frac{1}{d_\lambda} \text{Hf} [\tilde{S}] + C_o & \text{if } \ell(\lambda) = 2n-1 \text{ odd,} \end{cases}$$

where C_e and C_o are polynomials in odd Grassmannian variables of the order $\ell(\lambda)$, C_e starts with quadratic terms, while C_o starts with cubic ones. This follows from the consideration of (2.24) in Subsection 2.1.

4.2 Orthogonality

One can verify the equality, using (2.27)

$$\langle 0 | \Gamma(\mathbf{t}) | \lambda \rangle = \langle \lambda | \bar{\Gamma}(-\mathbf{t}) | 0 \rangle, \quad (4.9)$$

where

$$\bar{\Gamma}(\mathbf{t}) := e^{\bar{J}(\mathbf{t})}, \quad \bar{J}(\mathbf{t}) = \sum_{0 < i \in 1+2\mathbb{Z}} t_i J_{-i} + \sum_{0 < i \in \frac{1}{2}+\mathbb{Z}} t_i J_{-i} \quad (4.10)$$

(keep the order in the products of Grassmannian odd variables, which we label by semi-integer subscripts).

Thanks to (4.3) we can write

$$\langle 0 | \Gamma(\mathbf{t}) = \langle 0 | + \sum_{\lambda \in \text{OP}} C_\lambda(\mathbf{t}) \langle \lambda |. \quad (4.11)$$

On the other hand due to (4.9)

$$\bar{\Gamma}(\mathbf{t}) | 0 \rangle = | 0 \rangle + \sum_{\lambda \in \text{OP}} |\lambda \rangle C_\lambda(-\mathbf{t})$$

and to (2.18) we obtain

$$\langle 0 | \Gamma(\mathbf{t}) \bar{\Gamma}(\bar{\mathbf{t}}) | 0 \rangle = e^{-\frac{1}{2} \sum_{n=1,3,5,\dots} n t_n \bar{t}_n - \frac{1}{2} \sum_{m=0,1,2,\dots} (-1)^m (m+\frac{1}{2}) t_{m+\frac{1}{2}} \bar{t}_{m+\frac{1}{2}}},$$

where \mathbf{t} and $\bar{\mathbf{t}}$ are two independent sets of variables:

$$\mathbf{t} = (t_1, t_{\frac{1}{2}}; t_3, t_{\frac{3}{2}}; t_5, t_{\frac{5}{2}}; \dots), \quad \bar{\mathbf{t}} = (\bar{t}_1, \bar{t}_{\frac{1}{2}}; \bar{t}_3, \bar{t}_{\frac{3}{2}}; \bar{t}_5, \bar{t}_{\frac{5}{2}}; \dots).$$

On the other hand, due to (2.27), (4.10) and (4.3) we obtain the following analogue of Cauchy–Littlewood identity:

$$e^{-\frac{1}{2} \sum_{n=1,3,5,\dots} n t_n \bar{t}_n - \frac{1}{2} \sum_{m=1,2,3,\dots} (-1)^m (m+\frac{1}{2}) t_{m+\frac{1}{2}} \bar{t}_{m+\frac{1}{2}}} = \sum_{\lambda \in \text{OP}} C_\lambda(\mathbf{t}) C_\lambda(-\bar{\mathbf{t}}). \quad (4.12)$$

From the last equality we obtain

$$\Gamma(\mathbf{t}) = 1 + \sum_{\lambda \in \text{OP}} C_\lambda(\mathbf{t}) C_\lambda(\mathbf{J}), \quad \bar{\Gamma}(\mathbf{t}) = 1 + \sum_{\lambda \in \text{OP}} C_\lambda(\bar{\mathbf{J}}) C_\lambda(-\bar{\mathbf{t}}),$$

where

$$\begin{aligned} \mathbf{J} &:= \left(2J_1, 4J_{\frac{1}{2}}; \frac{2}{3}J_3, -\frac{4}{3}J_{\frac{3}{2}}; \frac{2}{5}J_5, \frac{4}{5}J_{\frac{5}{2}}; \frac{2}{7}J_7, -\frac{4}{7}J_{\frac{7}{2}}; \dots \right), \\ \bar{\mathbf{J}} &:= \left(-2J_{-1}, -4J_{-\frac{1}{2}}; -\frac{2}{3}J_{-3}, \frac{4}{3}J_{-\frac{3}{2}}; -\frac{2}{5}J_{-5}, -\frac{4}{5}J_{-\frac{5}{2}}; -\frac{2}{7}J_{-7}, \frac{4}{7}J_{-\frac{7}{2}}; \dots \right) \end{aligned}$$

and therefore

$$|\lambda\rangle = C_\lambda(\bar{\mathbf{J}}) |0\rangle, \quad \langle\lambda| = \langle 0| C_\lambda(\mathbf{J}). \quad (4.13)$$

Let $f(\mathbf{t})$ and $g(\mathbf{t})$ be series in the variables $\{t_i\}$. We introduce the following scalar product

$$\langle f, g \rangle := \langle 0| f(\mathbf{J}) g(\bar{\mathbf{J}}) |0\rangle. \quad (4.14)$$

In particular due to (2.18)

$$\langle t_i, t_j \rangle = \langle 0| J_i J_{-j} |0\rangle = -\frac{j}{2} (-1)^{[-j-\frac{1}{2}]} \delta_{i,j},$$

where $[a]$ denotes the integer part of a (i.e., $a = [a] + \epsilon$ where $0 \leq \epsilon \leq 1$ for $a < 0$).

It follows from (4.13) and (4.3) that the polynomials form an orthonormal basis in the scalar product (4.14):

$$\langle C_\lambda, C_\mu \rangle = \langle 0| C_\lambda(\mathbf{J}) C_\mu(\bar{\mathbf{J}}) |0\rangle = \delta_{\lambda,\mu}.$$

4.3 Polynomials C_λ in super Miwa variables

If we want to rewrite polynomials as symmetric functions symmetric with respect to the action of symmetric group on pairs (z_i, ζ_i) of super Miwa variables $\mathbf{z} = (z_1, \zeta_1, \dots, z_k, \zeta_k)$ (3.9), we present C_λ as

$$C_\lambda(\mathbf{z}) := C_\lambda(-[\mathbf{z}]) = \langle 0| :e^{\Theta(z_1, \zeta_1)} : \dots :e^{\Theta(z_k, \zeta_k)} : |\lambda\rangle D(\mathbf{z})^{-1}$$

in terms of super vertex operators (3.1). Here

$$D(\mathbf{z}) = D_k(\mathbf{z}) := \langle 0| :e^{\Theta(z_1, \zeta_1)} : \dots :e^{\Theta(z_k, \zeta_k)} : |0\rangle$$

is antisymmetric function with respect to the permutation of pairs (z_i, ζ_i) , $D(\mathbf{z})^{-1}$, see (3.10). Let us note that if $\ell(\lambda) > k$ the polynomial C_λ does not depend on the odd Grassmannian variables ζ_i , and vanish for $\ell(\lambda)$ odd. Example. For $k = 1$ we obtain

$$C_\lambda(z, \zeta) = \begin{cases} \frac{2}{d_\lambda} z^{-\frac{1}{2}|\lambda|} (2\ell(\lambda) - 1)!! & \text{if } \ell(\lambda) \text{ even,} \\ \frac{\zeta}{d_\lambda} z^{-\frac{1}{2}|\lambda| - \frac{1}{2}} & \text{if } \ell(\lambda) = 1, \\ 0 & \text{if } \ell(\lambda) > 1, \text{ odd,} \end{cases} \quad (4.15)$$

where we use the formula $s_{(n|m)}(x) = 2x^{n+m+1}$.

In Miwa variables the Cauchy–Littlewood identity (4.12) is written as

$$\sum_{\lambda \in \text{OP}} C_\lambda(\mathbf{z}) C_\lambda(-\bar{\mathbf{z}}) = \prod_{i,j} \frac{1 - z_i \bar{z}_j}{1 + z_i \bar{z}_j} \left(1 - \zeta_i \bar{\zeta}_j \frac{1 - z_i \bar{z}_j}{(1 + z_i \bar{z}_j)^2} \right).$$

Proof. Using (3.9), namely

$$t_{2n+1} = \frac{2}{(2n+1)} \sum_{i=1}^k \frac{1}{z_i^{2n+1}}, \quad t_{m+\frac{1}{2}} = 2 \sum_{i=1}^k \frac{\zeta_i}{z_i^m}$$

we obtain

$$\begin{aligned} e^{-\frac{1}{2} \sum_{m=1,2,3,\dots} (-1)^m (m+\frac{1}{2}) t_{m+\frac{1}{2}} \bar{t}_{m+\frac{1}{2}}} &= \prod_{i,j} \left(1 - 2\zeta_i \bar{\zeta}_j \sum_{m=1,2,3,\dots} (-1)^m \left(m + \frac{1}{2} \right) z_i^{-m} \bar{z}_j^{-m} \right) \\ &= \prod_{i,j} \left(1 + 2\zeta_i \bar{\zeta}_j D_{z_i} \sum_{m=1,2,3,\dots} (-1)^m z_i^{-m} \bar{z}_j^{-m} \right) \\ &= \prod_{i,j} \left(1 + 2\zeta_i \bar{\zeta}_j D_{z_i} \left(-\frac{1}{z_i \bar{z}_j} \frac{1}{1 + z_i^{-1} \bar{z}_j^{-1}} \right) \right) \\ &= \prod_{i,j} \left(1 - 2\zeta_i \bar{\zeta}_j D_{z_i} \left(\frac{1}{1 + z_i \bar{z}_j} \right) \right) = \prod_{i,j} \left(1 - \zeta_i \bar{\zeta}_j \frac{1 - z_i \bar{z}_j}{(1 + z_i \bar{z}_j)^2} \right), \end{aligned}$$

where $D_z = \frac{1}{2}(z\partial_z + \partial_z z)$ is the Euler operator. ■

4.4 Combinatorial meaning of $\mathbf{C}_\lambda(t) = \mathbf{C}_\lambda(1, 0, 0, \dots)$

Each partition with odd parts may be presented as $\lambda = (1^{m_1} 3^{m_3} 5^{m_5} \dots)$ where m_i is the multiplicity of the number i (that means that the partition λ contains the part equal to i m_i times). The length $\ell(\lambda)$ of the partition λ is equal to $\sum_{i=1,3,5,\dots} m_i$, the weight is $|\lambda| = \sum_{i=1,3,5,\dots} i m_i$.

Let us visualize this, in a similar way as in the papers [5, 18], as the one-dimensional semi-infinite lattice of sites (baskets) in our case numbered by odd positive integers. A basket number i ($i = 1, 3, 5, \dots$) contains m_i identical balls (and therefore the multiplicity m_i may be also called the *occupation number*). Nonequivalent distributions of balls is in one-to-one correspondence with partitions from the set OP. (The length of a partition is equal to the number of balls, the ratio of the weight of the partition and the length of the partition may be considered as the location of the mass center of the balls).

Let us consider the following discrete time random process describing the creation of $\lambda \in \text{OP}$ or, the same, of ball configurations. It starts with a given partition, say, μ . (The case where $\mu = 0$ describes the configuration where all baskets are empty at time $\tau = 0$). At each discrete time instant $\tau = 1, 2, 3, \dots$ one of the following two possible events occurs with equal probability (A) either two balls are created in the leftmost basket (basket number 1), or (B) a ball chosen at random in any of baskets, say, in basket number i , is moved to the nearest basket to the right (to the basket number $i+2$). It is clear that at each time step the weight of the related partition increases: $|\lambda| \rightarrow |\lambda| + 2$, thus $|\lambda| = 2\tau$. A problem is to find a number of ways to create a given distribution λ of the balls in baskets along the process described above in $\tau = \frac{1}{2}|\lambda|$ steps. We denote this number $\mathbf{N}_{\mu \rightarrow \lambda}$.

Then we state that

$$\mathbf{C}_\lambda(1, 0, 0, 0, \dots) = \left(\frac{1}{2} \right)^{\frac{1}{2}\ell(\lambda)} \frac{1}{(\frac{1}{2}|\lambda|)!} \frac{1}{d_\lambda} \mathbf{N}_{0 \rightarrow \lambda}$$

The proof follows from

$$e^{J^{-1}}|0\rangle = |0\rangle + \sum_{\lambda \in \text{OP}_e} |\lambda\rangle \mathbf{C}_\lambda(-1, 0, 0, 0, \dots) = \sum_{\tau=0,1,2,\dots} \frac{1}{\tau!} (J_{-1})^\tau |0\rangle$$

and from the detailed consideration of the action of

$$J_{-1} = -\frac{1}{2}\phi_{\frac{1}{2}}^2 + \phi_{-\frac{1}{2}}\phi_{\frac{3}{2}} - \phi_{-\frac{3}{2}}\phi_{\frac{5}{2}} + \dots$$

on basis Fock vectors (4.1) (see Appendix B), and from (4.8).

Now we write down the following formula for $C_\lambda(1, 0, 0, \dots)$ obtained from (4.7):

$$C_\lambda(1, 0, 0, \dots) = \frac{1}{d_\lambda} \prod_i \frac{1}{n_i! n_j!} \text{Hf} \left[\frac{1}{n_i + n_j + 1} \right],$$

where n_i and λ_i are related by (4.2), which yields

$$N_{0 \rightarrow \lambda} = 2^{\frac{1}{2}\ell(\lambda)} \left(\frac{1}{2}|\lambda| \right)! \left(\prod_{i=1,3,5,\dots} \frac{1}{m_i!} \right) \text{Hf} \left[\frac{1}{n_i! n_j! (n_i + n_j + 1)} \right],$$

where we recall that m_i and n_i we defined above: $\lambda = (1^{m_1} 3^{m_3} 5^{m_5} \dots)$ and $\lambda = (\lambda_1, \dots, \lambda_{2k})$ where $\ell(\lambda) = 2k$ is the length of the partition, $\lambda_i = 2n_i + 1$.

For instance, take $\lambda = (1^{2k})$. Obviously there is only one way to create this configuration, $N_{0 \rightarrow (1^{2k})} = 1$. Indeed $\ell(\lambda) = |\lambda| = 2k$, $n_i = 0$, $i = 1, \dots, 2k$, and we obtain well-known identity

$$1 = N_{0 \rightarrow (1^{2k})} = 2^k k! \frac{1}{(2k)!} (2k-1)!!.$$

At last let us mention that $C_\lambda(\mathbf{t})$ where $t_i = 0$ except $i = \frac{1}{2}$, 1 may be related to the numbers $N_{(1) \rightarrow \lambda}$, where $\lambda \in \text{OP}_o$.

4.5 Polynomials $C_{\lambda/\mu}$

We define skew polynomials $C_{\lambda/\mu}(\mathbf{t})$ as follows (cf. (2.27) and (4.10)):

$$C_{\lambda/\mu}(\mathbf{t}) := \langle \mu | \Gamma(\mathbf{t}) | \lambda \rangle = \langle \lambda | \bar{\Gamma}(-\mathbf{t}) | \mu \rangle.$$

From this definition $C_{\lambda/\mu}$ vanishes unless $\mu \subseteq \lambda$. The same may be written as

$$\Gamma(\mathbf{t}) = \sum_{\lambda, \mu} C_{\lambda/\mu}(\mathbf{t}) |\mu\rangle \langle \lambda|, \quad \bar{\Gamma}(\mathbf{t}) = \sum_{\lambda, \mu} C_{\lambda/\mu}(-\mathbf{t}) |\lambda\rangle \langle \mu|.$$

If $\mathbf{t} = \mathbf{t}' + \mathbf{t}''$ then $\Gamma(\mathbf{t}) = \Gamma(\mathbf{t}')\Gamma(\mathbf{t}'')$ by inserting the unity operator $\sum_{\lambda \in P} |\lambda\rangle \langle \lambda|$ between $\Gamma(\mathbf{t}')$ and $\Gamma(\mathbf{t}'')$ we obtain

$$C_\lambda(\mathbf{t}) = \sum_{\mu} C_{\lambda/\mu}(\mathbf{t}') C_\mu(\mathbf{t}'').$$

This property is quite similar to the property of the Schur functions (see (5.9) in I of [11]).

One may relate $C_{\lambda/\mu}$ to the numbers $N_{\mu \rightarrow \lambda}$ described in the previous subsection.

4.6 CKP tau function and polynomials C_λ

First of all we note that $C_\lambda(\mathbf{t})$ is not a solution of the Hirota bilinear equations, and, therefore is not a CKP tau function. However, due to (4.11) CKP tau functions (3.6) are series in $C_\lambda(\mathbf{t})$ as follows

$$\tau(\mathbf{t}) = \sum_{\lambda \in \text{OP}} g_\lambda C_\lambda(\mathbf{t}),$$

where

$$g_\lambda = \langle \lambda | g | 0 \rangle.$$

Example 1. Take $g = \exp \sum_{i>0} e^{-2U_i} \phi_i^2$. Then

$$\tau(\mathbf{t}) = \sum_{\lambda} e^{-U_{\lambda}} \mathbf{C}_{\lambda}(\mathbf{t}) d_{\lambda} \prod_{i=1,3,5,\dots} \frac{1}{k_i!}, \quad (4.16)$$

where the sum ranges over all $\lambda \in \text{OP}$ whose parts have even multiplicities, $m_i =: 2k_i$, i.e., of the form $\lambda = (1^{2k_1} 3^{2k_3} 5^{2k_5} \dots)$. The numbers U_{λ} are defined as

$$U_{\lambda} := \sum_{i=1}^k U_{\frac{1}{2}\lambda_i}. \quad (4.17)$$

The right-hand side of (4.16) may be compared to sums over partitions in [6] and in [14] dealing with tau functions of neutral and charged BKP hierarchies, respectively.

Example 2. Given a symmetric matrix A and a partition λ introduce numbers A_{λ} according to the formula

$$e^{\sum_{n,m>0} x_n A_{nm} x_m} = \sum_{\lambda \in P_{\text{ev}}} A_{\lambda} x_{\lambda},$$

where $\lambda = (\lambda_1, \dots, \lambda_k)$ is a partition, and $x_{\lambda} = x_{\lambda_1} \cdots x_{\lambda_k}$, and where P_{ev} is a set of all partitions with even number of parts.

Then taking

$$e^{\sum_{n,m>0} \phi_{\frac{n}{2}} A_{nm} \phi_{\frac{m}{2}}},$$

which is an exponential of a quadratic form of creation operators, we obtain

$$\tau(\mathbf{t}) = \langle 0 | \Gamma(\mathbf{t}) g | 0 \rangle = \sum_{\lambda \in P_{\text{ev}}} A_{\lambda} \mathbf{C}_{\lambda}(\mathbf{t}).$$

Next we may write the CKP tau function as a double series over partitions

$$\tau(\mathbf{t}, \bar{t}) = \sum_{\lambda, \mu \in \text{OP}} g_{\lambda, \mu} \mathbf{C}_{\lambda}(\mathbf{t}) \mathbf{C}_{\mu}(\bar{t}),$$

where

$$g_{\lambda, \mu} = \langle \lambda | g | \mu \rangle.$$

Example 3.

$$\tau(\mathbf{t}, U, \bar{t}) = \langle 0 | \Gamma(\mathbf{t}) \mathbb{T}(U) \bar{\Gamma}(\bar{t}) | 0 \rangle = 1 + \sum_{\lambda \in \text{OP}} e^{-U_{\lambda}} \mathbf{C}_{\lambda}(\mathbf{t}) \mathbf{C}_{\lambda}(-\bar{t}),$$

where $U = (U_{\frac{1}{2}}, U_{\frac{3}{2}}, U_{\frac{5}{2}}, \dots)$ is a set of constants, and

$$\mathbb{T}(U) := \exp \sum_{0 < i \in \frac{1}{2} + \mathbb{Z}} (-1)^{i+\frac{1}{2}} U_i \varphi_i \varphi_{-i},$$

where the numbers U_{λ} are defined by (4.17).

This example may be compared with the results of [15] and [13] devoted to TL and neutral BKP hypergeometric tau functions.

We can specify \bar{t} in such a way that $\bar{t}_i = \frac{2}{i} z^i$, $i = 1, 3, 5, \dots$, and use (4.15). We thus obtain

$$\tau(\mathbf{t}, U, \bar{t}(z)) = 1 + \sum_{\lambda \in \text{OP}_{\text{ev}}} f_{\lambda} e^{-U_{\lambda}} \mathbf{C}_{\lambda}(\mathbf{t}), \quad f_{\lambda} := (-1)^{\frac{1}{2}(|\lambda| + \ell(\lambda))} \frac{2}{d_{\lambda}} z^{\frac{1}{2}|\lambda|} (2\ell(\lambda) - 1)!!.$$

A Examples of CKP tau functions

(I) One-soliton tau function is

$$\langle 0 | \Gamma(t) e^{\frac{a}{2} \phi(p) \phi(q)} | 0 \rangle = \left(1 - \frac{a}{p+q} e^{\sum (p^n + q^n) t_n} \right)^{-\frac{1}{2}}.$$

(II) Another example is as follows

$$\langle 0 | \Gamma(t) e^{a \phi_{\frac{1}{2}}^2} | 0 \rangle = (1 - at_1)^{-\frac{1}{2}},$$

which may be also viewed as a solution of a heat equation

$$\frac{\partial f(z_{\frac{1}{2}}, t_1)}{\partial t_1} = \frac{\partial^2 f(z_{\frac{1}{2}}, t_1)}{\partial z_{\frac{1}{2}}^2}, \quad f(z_{\frac{1}{2}}, 0) = \exp a z_{\frac{1}{2}}^2,$$

which is

$$f(z_{\frac{1}{2}}, t_1) = \exp t_1 \partial_{\frac{1}{2}}^2 \cdot \exp a z_{\frac{1}{2}}^2 = \frac{1}{\sqrt{1 - at_1}} \exp \frac{a z_{\frac{1}{2}}^2}{1 - at_1}$$

in the origin.

B A realization of the algebra of free bosons (2.1)–(2.3)

The simplest way to understand the action of J_{-1} on a basis Fock vector may be as follows. Consider the realization of free boson algebra (2.1) via differentiation operators in auxiliary variables z_i :

$$\phi_{m+\frac{1}{2}} = z_{m+\frac{1}{2}}, \quad \phi_{-m-\frac{1}{2}} = (-)^m \partial_{m+\frac{1}{2}}, \quad m \geq 0,$$

while the Fock space F may be viewed as polynomial functions in the auxiliary variables z_i

$$| (1^{m_1} 3^{m_3} 5^{m_5} \dots) \rangle d_\lambda = z_{\frac{1}{2}}^{m_1} z_{\frac{3}{2}}^{m_3} z_{\frac{5}{2}}^{m_5} \dots.$$

The action of J_{-1} on a basis Fock vector yields the following

$$\begin{aligned} & \left(-\frac{1}{2} z_{\frac{1}{2}}^2 + z_{\frac{3}{2}} \partial_{\frac{1}{2}} + z_{\frac{5}{2}} \partial_{\frac{3}{2}} + \dots \right) z_{\frac{1}{2}}^{m_1} z_{\frac{3}{2}}^{m_3} z_{\frac{5}{2}}^{m_5} \dots \\ &= -\frac{1}{2} z_{\frac{1}{2}}^{m_1+2} z_{\frac{3}{2}}^{m_3} z_{\frac{5}{2}}^{m_5} \dots + m_1 z_{\frac{1}{2}}^{m_1-1} z_{\frac{3}{2}}^{m_3+1} z_{\frac{5}{2}}^{m_5} \dots + m_3 z_{\frac{1}{2}}^{m_1} z_{\frac{3}{2}}^{m_3-1} z_{\frac{5}{2}}^{m_5+1} \dots + \dots \end{aligned}$$

As we see the action of J_{-1} on a basis vector results in linear combination of basis vectors. The first term in the right-hand side of the last equality may be related to the event (A) (creation of a pair of balls in the basket 1. This event is accomplished by the multiplication by the factor $-\frac{1}{2}$. Thus we get the general prefactor $(-1)^{\ell(\lambda)}$ for all configurations with $\ell(\lambda)$ number of the balls. Other terms in the right-hand side may be related to the event (B) where a chosen ball is moved to the right neighbor basket, each possibility has the weight 1. The factors m_i describe the fact that we can chose any of m_i balls in the basket i to move them to the right.

C Bosonic KP tau function

Here we shall briefly describe a bosonic KP tau function and its symplectic reduction to the bosonic CKP tau function.

Consider the following bosonic operators

$$[p_i, q_j] = \delta_{i,j}, \quad i, j \in \frac{1}{2} + \mathbb{Z}.$$

The right and left Fock spaces will be defined via

$$p_i|0\rangle = q_{-i}|0\rangle = 0 = \langle 0|q_i = \langle 0|p_{-i}, \quad i < 0.$$

Let us define the expression $p_i q_j - \langle 0|p_i q_j|0\rangle$ by $E_{ij} = :p_i q_j:$. These E_{ij} may be considered as generators of $gl(\infty)$ algebra with a negative level.

Let us notice that for

$$g := \exp \sum_{i,j \in \frac{1}{2} + \mathbb{Z}} A_{ij} :p_i q_j: \quad (\text{C.1})$$

we have

$$\sum_{i \in \frac{1}{2} + \mathbb{Z}} g p_i \otimes g q_i = \sum_{i \in \frac{1}{2} + \mathbb{Z}} p_i g \otimes q_i g.$$

Bosonic KP tau function may be defined as

$$\tau^{KP}(\mathbf{t}) := \langle 0|\Gamma(\mathbf{t})g|0\rangle, \quad \Gamma(\mathbf{t}) := \exp \sum_{n=1}^{\infty} t_n \sum_{k \in \mathbb{Z}} p_k q_{k+n},$$

where $\mathbf{t} = (t_1, t_2, \dots)$ is a set of higher times. This tau function may be considered as a particular case of a supersymmetric KP tau function [9].

If we ask g of (C.1) to be invariant under $Sp(\infty)$ group then g may be expressed in terms of the so-called symplectic bosons.

Symplectic bosons. Now between the bosons

$$\phi_i = \frac{1}{\sqrt{2}} \left(p_i - (-1)^{i+\frac{1}{2}} q_{-i-1} \right), \quad \hat{\phi}_i = \frac{1}{\sqrt{2}} \left(p_i + (-1)^{i+\frac{1}{2}} q_{-i-1} \right)$$

there are the following relations, cf. (2.1), (2.2):

$$[\phi_i, \phi_j] = [\hat{\phi}_j, \hat{\phi}_i] = (-)^{j-\frac{1}{2}} \delta_{i,-j} \quad \text{and} \quad [\phi_i, \hat{\phi}_j] = 0.$$

It is known (see [9]) that quadratic expressions $Z_{ij} = : \phi_i \phi_j :$ ordered via (2.6) may be considered as a realization for the generators of the c_∞ algebra with a negative level (this fact may be verified with the help of (2.1) and (2.6)). The same is true for $\hat{Z}_{ij} = : \hat{\phi}_i \hat{\phi}_j :$.

It is straightforward to show that if g of (C.1) is invariant under the $Sp(\infty)$ group then $g = g_1 \hat{g}_1$, where $g_1 = \exp \left(\sum_{i,j \in \frac{1}{2} + \mathbb{Z}} A_{ij} : \phi_i \phi_j : \right)$ and $\hat{g}_1 = \exp \left(\sum_{i,j \in \frac{1}{2} + \mathbb{Z}} A_{ij} : \hat{\phi}_i \hat{\phi}_j : \right)$; then, if all

$t_{2n-1} = 0$, the KP bosonic tau function may be factorized: $\tau^{KP} = (\tau^{CKP})^{-2}$, where τ^{CKP} is of form (2.31), $\alpha = 0$ (see also (3.6)).

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